

Searches for Charged Higgs at CDF II

Ricardo Eusebi

Fermi National Accelerator Laboratory, CDF

(on behalf of the CDF collaboration)

Outline :

- × *Introduction*
- × *Single Channel search*
- × *Multi-channel search*
 - × *Results for Tevatron*
 - × *Expectation for LHC*
- × *Ongoing analyses*
- × *Summary*



Introduction

Electroweak Symmetry Breaking

→ Standard Model (SM) : 1 Higgs doublet

- EWSB → One Higgs boson, $h^{0(\text{SM})}$
- Decays to bb , $\tau\tau$, etc.

→ Natural Next step : Models with 2 Higgs Doublets (2HDM)

- EWSB → 5 Higgs bosons (h^0, H^0, A, H^\pm)
- $h^0, H^0 \rightarrow bb, \tau\tau, gg, W^+W^-, ZZ, cc$
- $A \rightarrow bb, \tau\tau, gg, Zh^0, tt$
- $H^\pm \rightarrow tb, \tau\nu, cs, W^\pm h^0, W^\pm A$

↑ **direct evidence of physics beyond SM**

→ Top

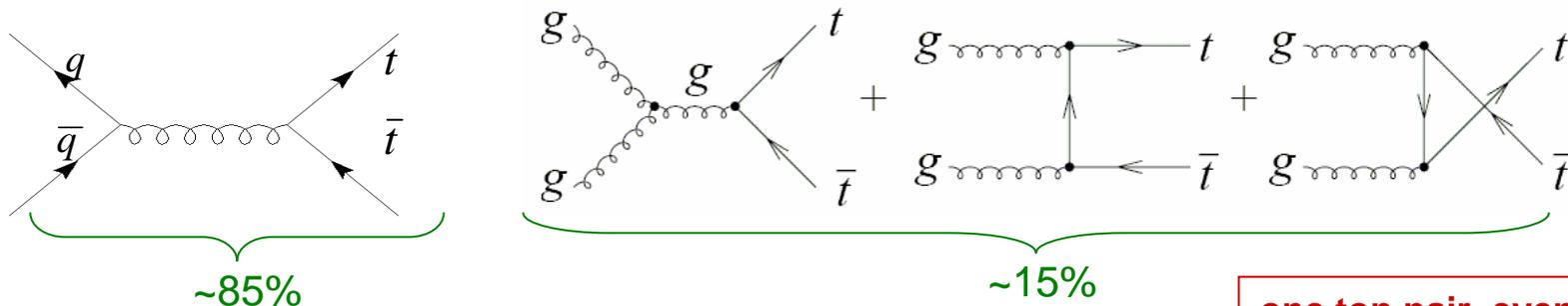
- Large mass suggest it plays an important role
- Fermion to which coupling to Higgs is most important, $y_t = M_t/v \approx 1$.

Top and Higgs bosons are the main characters in EWSB

→ **What can be said about Top and charged Higgs production?**

Top quark production and decay

Produced in pairs via the strong interactions



$$\sigma(\bar{p}p \rightarrow t\bar{t} @ M_{top} = 175 GeV, \sqrt{s} = 1.96 TeV) \approx 6.7 \text{ pb}$$

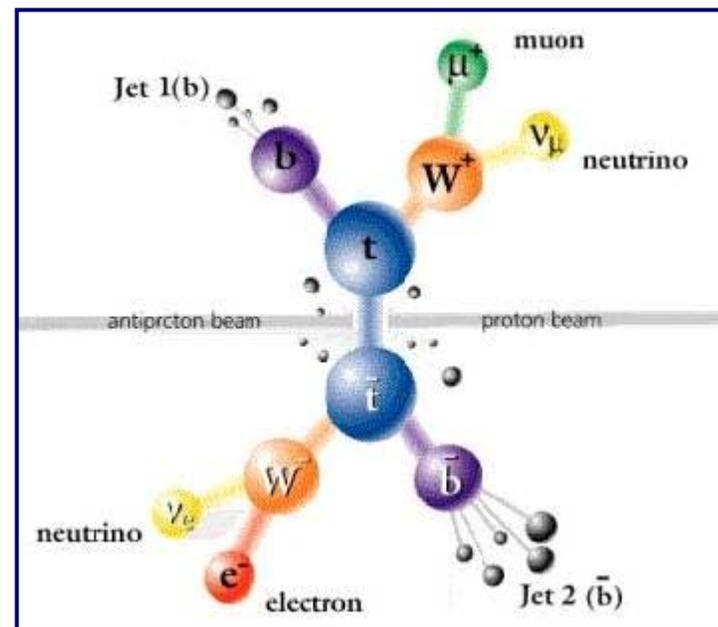
one top pair event every 10^{10} inelastic collisions

Other mechanisms of production possible, but not yet observed.

Decay:

In the SM $BR(t \rightarrow W^+b) > 0.99$ @95 C.L.

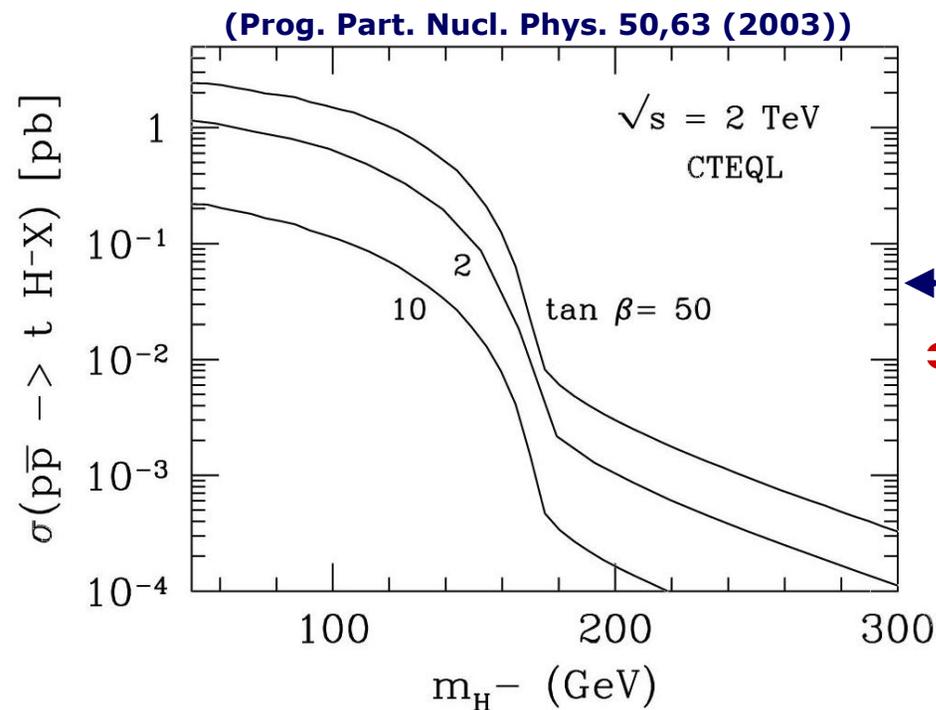
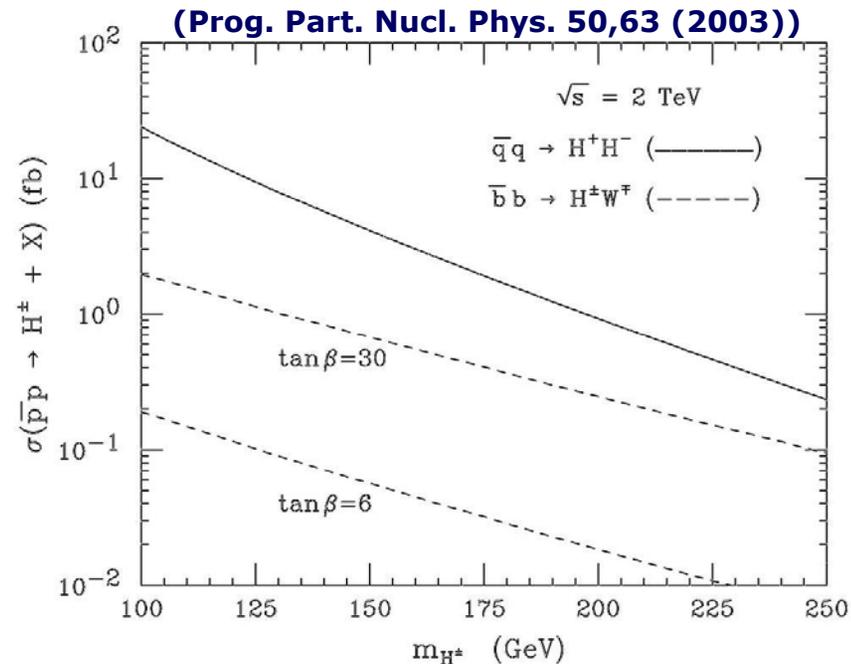
$$BR(t \rightarrow H^+b) = ??$$



Charged Higgs Production

Direct C.H. Production

- $\sqrt{s}=2$ TeV : $q\bar{q} \rightarrow H^+H^-$
- Very small production rate
- Signature hard to distinguish



Indirect C.H. Production

- $\sqrt{s}=2$ TeV : top associated
 - If $m_{\text{top}} > m_H + m_b$ from $t\bar{t}$ decays
 - If $m_{\text{top}} < m_H$ with associated top
- **Maybe Large production rates !**
 - **compare with ~ 6.7 pb**
- **Clean signature !**

Single-channel search (335 pb⁻¹)

Anomalous tau production in tt events

(tt, $t \rightarrow H^+ b$, $H^+ \rightarrow \underline{\tau} \nu$)

Event Selection

- ➔ **One high p_T electron or muon (trigger)**
 - ➔ $E_T > 20$ GeV
- ➔ **At least two jets.**
 - ➔ Leading jet $E_T > 25$ GeV, other jet $E_T > 15$ GeV
 - ➔ Cone radius $R = (\eta^2 + \phi^2)^{1/2} > 0.4$
- ➔ **At least one jet identified as a b-jet.**
 - ➔ Efficiency in $t\bar{t}$: 48 %
 - ➔ Displaced vertex
- ➔ **Missing transverse energy, $E_T > 20$ GeV.**
- ➔ **Scalar sum of all objects, $H_T > 205$ GeV.**
- ➔ **Exactly one hadronically decaying Tau lepton.**
 - ➔ Tau-jet likelihood > 0.65

Tau-jet likelihood

⇒ Hadronic Tau candidates distinguished in 4 types

Tau Type	No. of tracks	No. of π^0 s
Type1	1	0
Type2	1	≥ 1
Type3	3	≥ 0
Type4	2	≥ 0

⇒ One Likelihood per type. Based on variables characterizing a tau

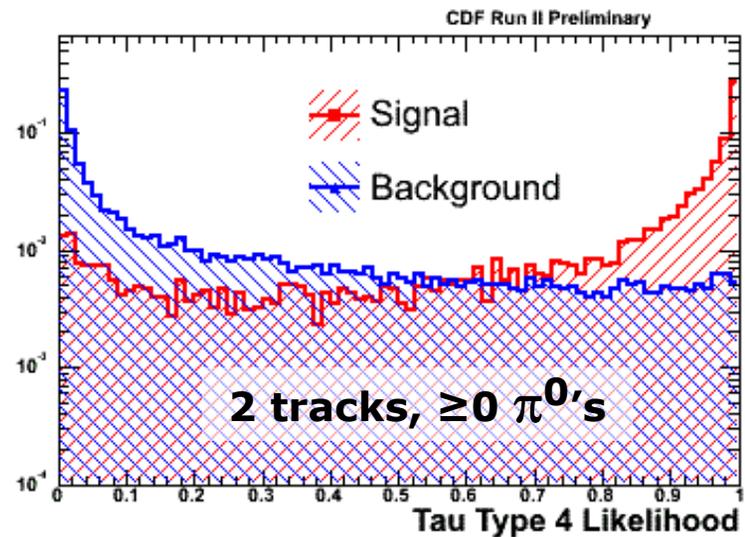
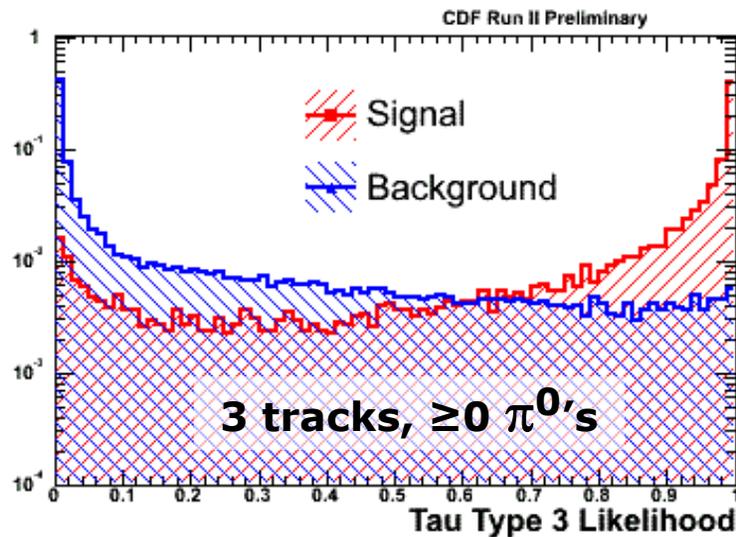
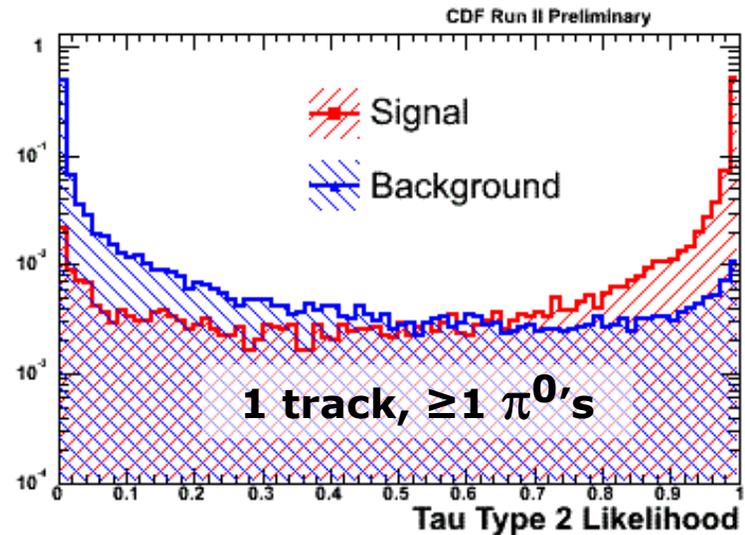
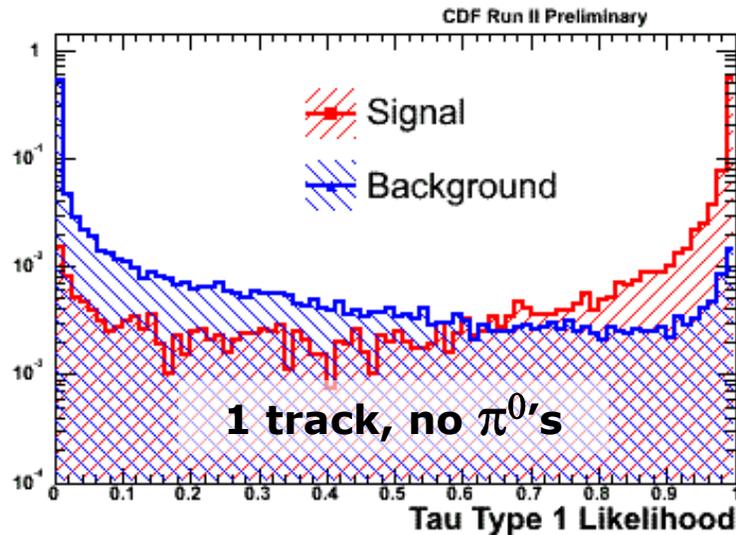
- ⇒ Tracking and calorimeter isolation
- ⇒ Ratios of track P_T
- ⇒ Maximum d_0 significance

$$\mathcal{L} = \frac{\prod_i^N \mathbf{P}_S^i(x_i)}{\prod_i^N \mathbf{P}_S^i(x_i) + \prod_i^N \mathbf{P}_B^i(x_i)}$$

where \mathbf{P}_S (\mathbf{P}_B) is the signal (background) probability for variable i .

Tau-jet likelihood fitter

Signal : $t\bar{t}$ Monte Carlo.
Background: jet-triggered data



Results

→ Systematic uncertainties

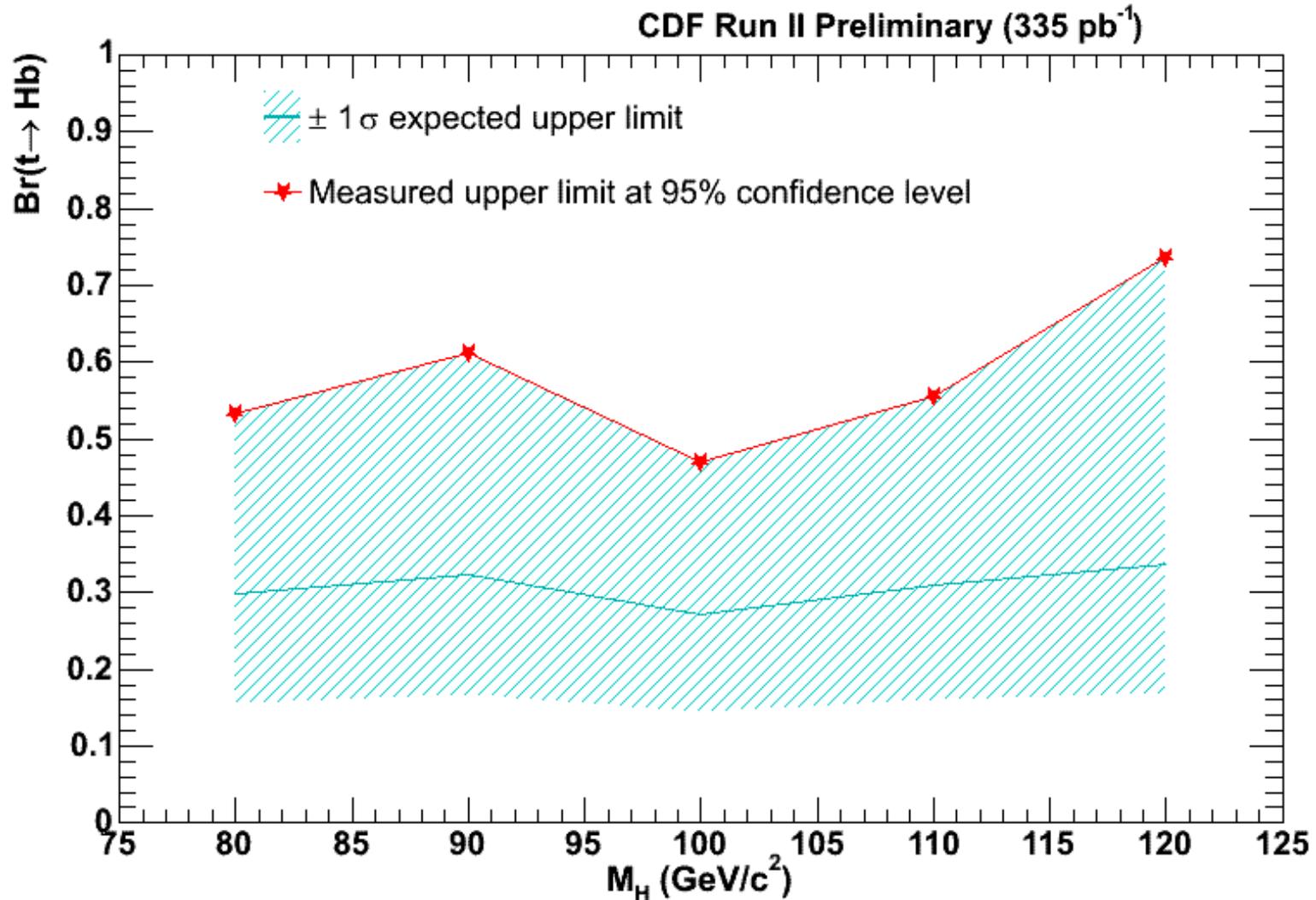
Source	Relative Error
jet \rightarrow τ misidentification	20%
Tau identification	15%
b -jet identification	7%
jet \rightarrow b -jet misidentification	11%
Jet energy scale	5%
Lepton identification	2%

→ Results

	Electron, Tau	Muon, Tau	All
$t\bar{t} \rightarrow \tau$	1.22 ± 0.22	0.85 ± 0.15	2.07 ± 0.37
fake τ, b -jet	0.65 ± 0.14	1.10 ± 0.22	1.74 ± 0.36
Other	0.03 ± 0.03	0.02 ± 0.02	0.06 ± 0.06
Total	1.90 ± 0.26	1.97 ± 0.27	3.88 ± 0.52
Data	4	2	6
Prob.	0.13	0.58	0.20

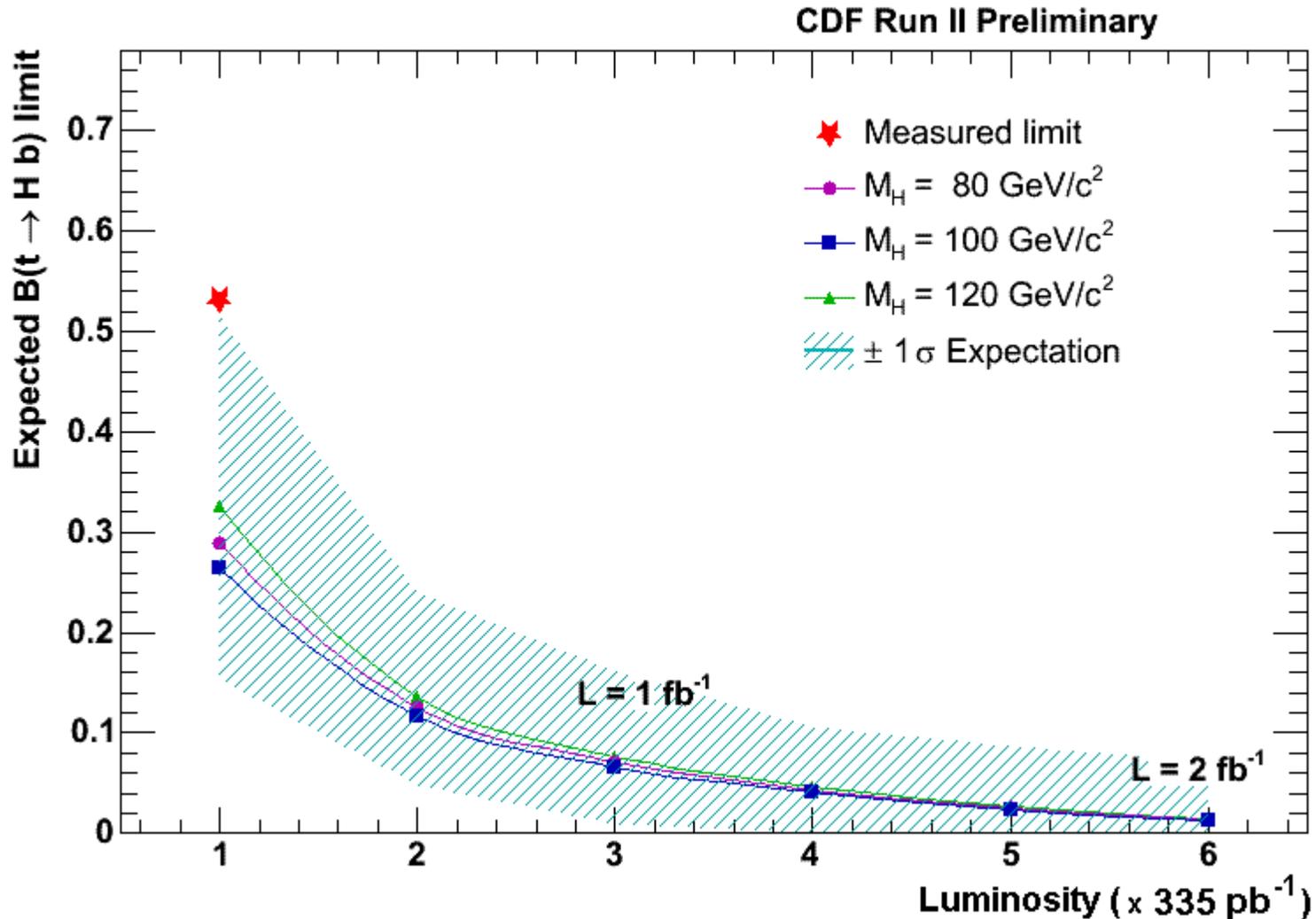
Results

→ In terms of $\text{BR}(t \rightarrow H^+ b)$ Vs. m_{H^+} , assuming $\sigma_{tt} = 6.7 \text{ pb}$.



Expected limit Vs. Luminosity

Expected limits on $BR(t \rightarrow H^+ b)$ Vs. Lum.



Multi-channel search (200 pb⁻¹)

PRL 96, 042003 (2006)

tt, $t \rightarrow H^+ b$, $H^+ \rightarrow c \underline{s}$, $t^ \underline{b}$, Wh^0 , τ ν*

Charged Higgs Decays Considered

- ⇒ Assume that top decays either to W^+b or H^+b
 - ⇒ $t \rightarrow W^+b$ $BR(t \rightarrow W^+b) \equiv 1 - BR(t \rightarrow H^+b)$
 - ⇒ $t \rightarrow H^+b$
- ⇒ Assume that the Higgs decays only as follows :
 - ⇒ $H^+ \rightarrow c\bar{s}$ $BR(H^+ \rightarrow \tau\nu) \equiv 1 - BR(H^+ \rightarrow c\bar{s})$
 - ⇒ $H^+ \rightarrow \tau\nu$ $- BR(H^+ \rightarrow t^*b)$
 - ⇒ $H^+ \rightarrow t^*b$ $- BR(H^+ \rightarrow W^+h^0)$
 - ⇒ $H^+ \rightarrow W^+h^0$
- ⇒ We further consider the h^0 decays to $b\bar{b}$
 - ⇒ $h^0 \rightarrow b\bar{b}$ $BR(h^0 \rightarrow b\bar{b})$
- ⇒ **Summary** : For each top quark we consider 5 possible decays modes

B₁. $t \rightarrow W^+b$	
B₂. $t \rightarrow H^+b \rightarrow c\bar{s}b$	B₃. $t \rightarrow H^+b \rightarrow \tau\nu b$
B₄. $t \rightarrow H^+b \rightarrow t^*b\bar{b}$	B₅. $t \rightarrow H^+b \rightarrow W^+h^0b \rightarrow W^+bb\bar{b}$

The BR to each (B_i) can be predicted from these 5 indep. BR's
The Narrow Width Approximation (NWA) is implicit.

Search Strategy

Look at the relative rates of events in different $t\bar{t}$ decay channels

➔ **Take advantage of existent cross section analyses**

- ➔ Lepton+Jets (1, and 2 or more tags) Phys. Rev. **D71**,052003
- ➔ Lepton+Tau Phys. Lett. **B639**, 172 (2006)
- ➔ Dilepton PRL. **93**, 142001 (2004)

➔ **If H^\pm is present, and $t \rightarrow H^+ b$, what do we expect ?**

- ➔ If **$BR(H^+ \rightarrow \tau \nu) \sim 1$** :
 - ➔ Lepton+tau will see an excess w.r.t. SM expectations.
 - ➔ Dilepton and Lepton+jets will show a deficit.
- ➔ If **$BR(H^+ \rightarrow c s) \sim 1$** :
 - ➔ All channels would see a deficit. All hadronic channel will see an excess.
- ➔ Similar considerations for other H^+ decays.

Analysis Technique

Model : Top
and H^\pm
production



#events in
search channels



Comparison
to data



Limits on
model

Results

Presented for three types of model:

- ➔ The MSSM
- ➔ The tauonic Higgs model
- ➔ The Higgs BR independent model.

Results: MSSM , Choice of Benchmarks

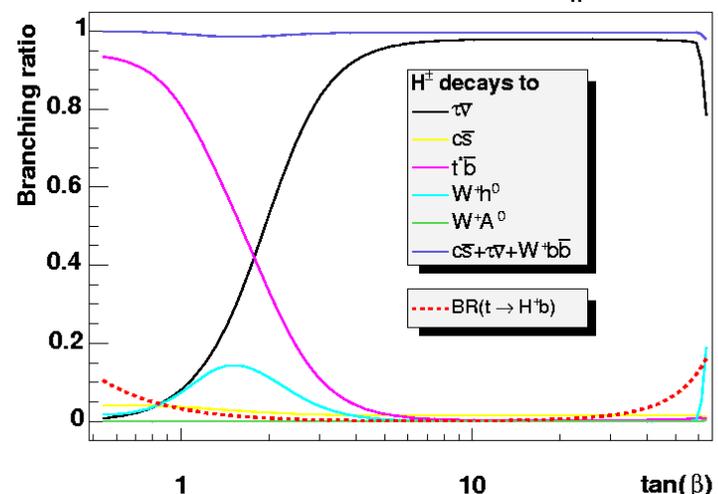
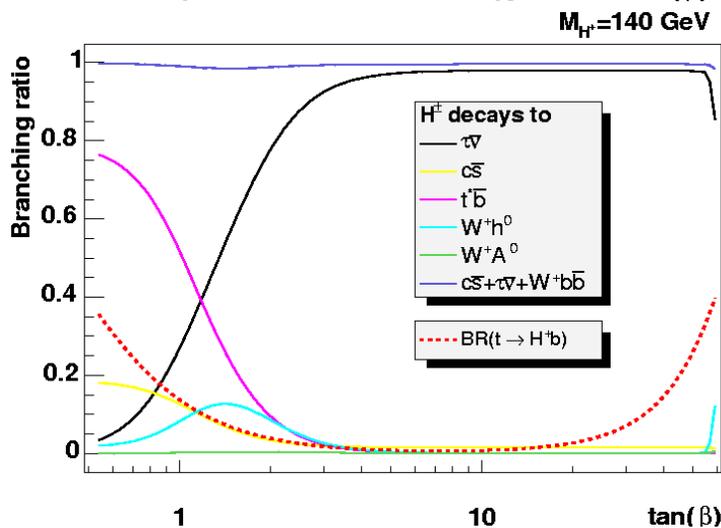
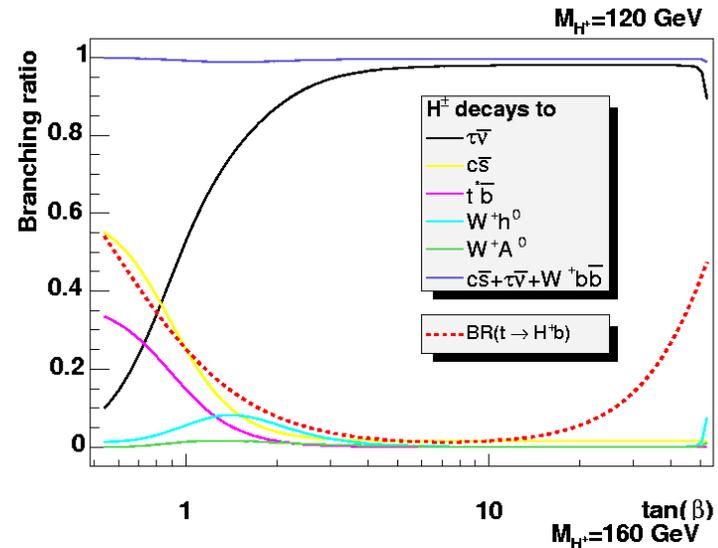
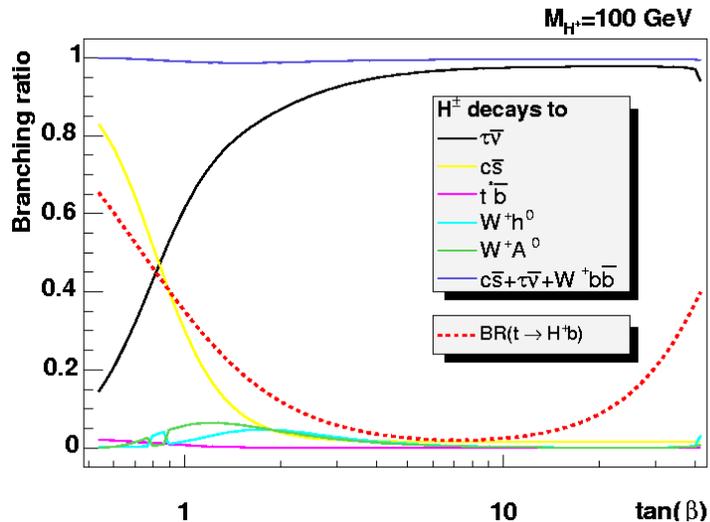
Using Phenomenological MSSM; 14 params for the 3rd generation.

	μ (GeV)	A_t (GeV)	Rest of parameters (GeV)
B1	-500	2000	$M_2=M_3=M_Q=M_U=M_D=1$ TeV $M_1=0.4978*M_2$, $M_L=M_E= 1$ TeV $A_b=A_t$, $A_{\tau}=500$ GeV
B2	-500	-500	
B3	500	500	
B4	500	2800	
B5(Minimal)	-200	$\mu/\tan(\beta)$	$M_Q=M_U=M_D=1$ TeV, $M_2=M_3=200$ $M_E=M_L=M_Q$, $A_t=A_b$, $A_{\tau}=500$
B6(Maximal)	-200	$2450\text{GeV} + \mu/\tan(\beta)$	

- ➔ **B1 and B2 value of $\mu=-500$ GeV, large $BR(t \rightarrow H^+b)$ at large $\tan(b)$**
 - ➔ Difference is A_t , that is chosen so as to maximize (B1) and minimize (B2) the mass of the h^0 in the $\tan(b) \sim 1$ region.
- ➔ **B3 and B4 value of $\mu=+500$ GeV, small $BR(t \rightarrow H^+b)$ at large $\tan(b)$**
 - ➔ Difference is A_t , that is chosen so as to maximize (B4) and minimize (B3) the mass of the h^0 in the $\tan(b) \sim 1$ region.
- ➔ **B5 and B6 are the minimal and maximal stop mixing scenarios used at LEP. They minimize and maximize the mass of the h^0 at every point in $\tan(b)$.**

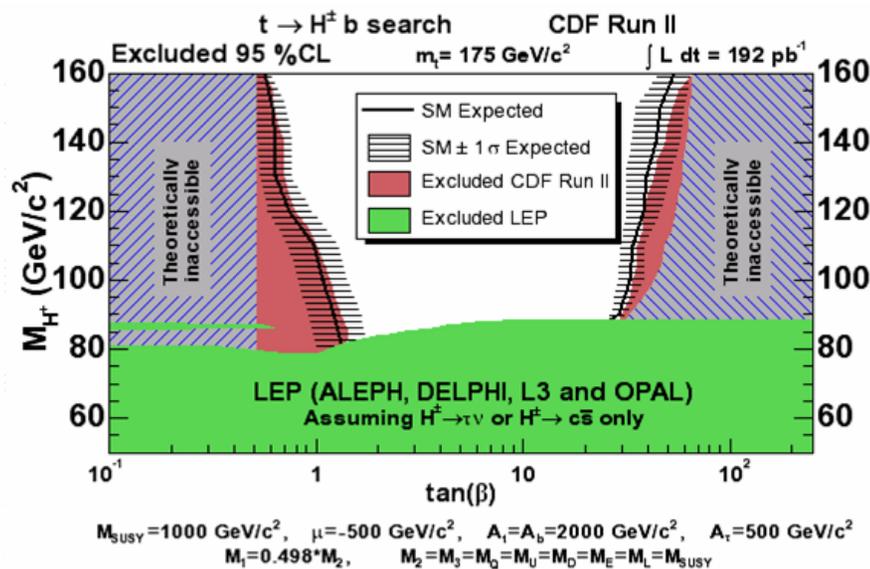
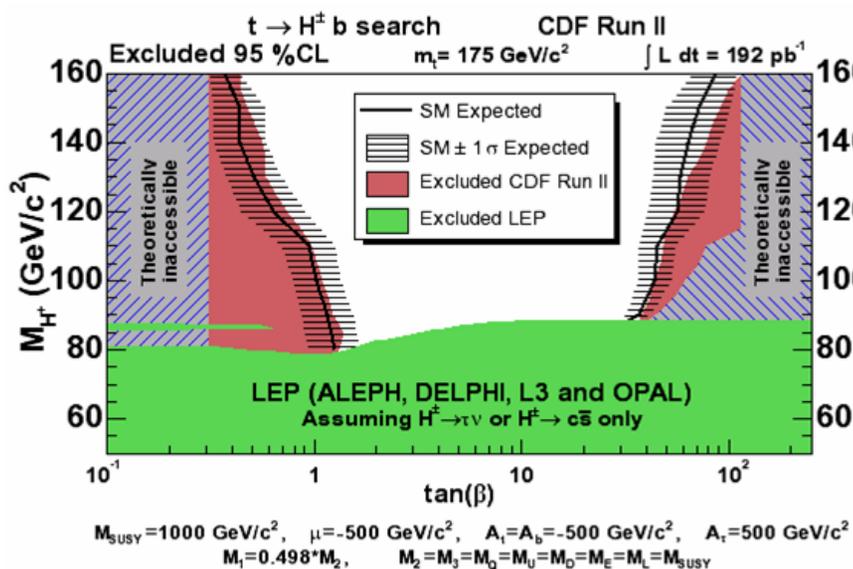
Results: MSSM, BR Predictions

➔ CPsuperH (hep-ph/0307373) predicts the Higgs BR's. For Bench #1 :

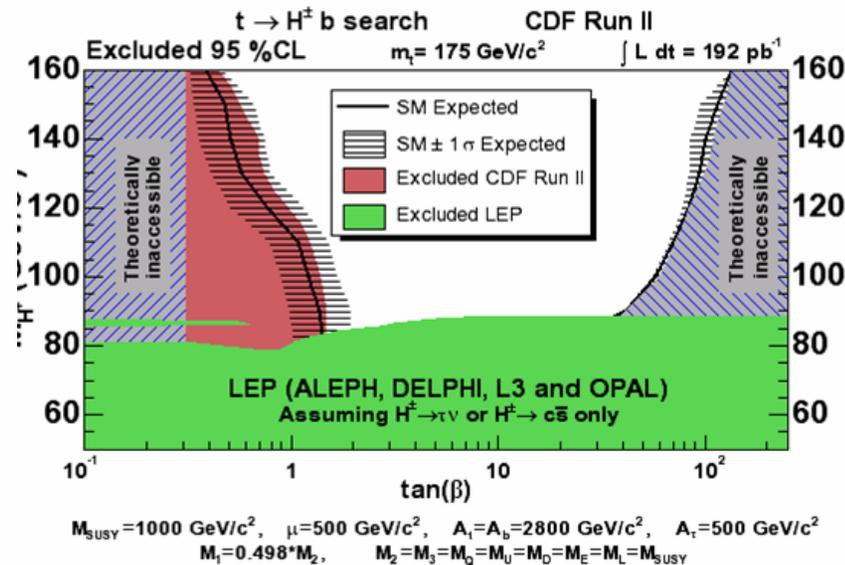
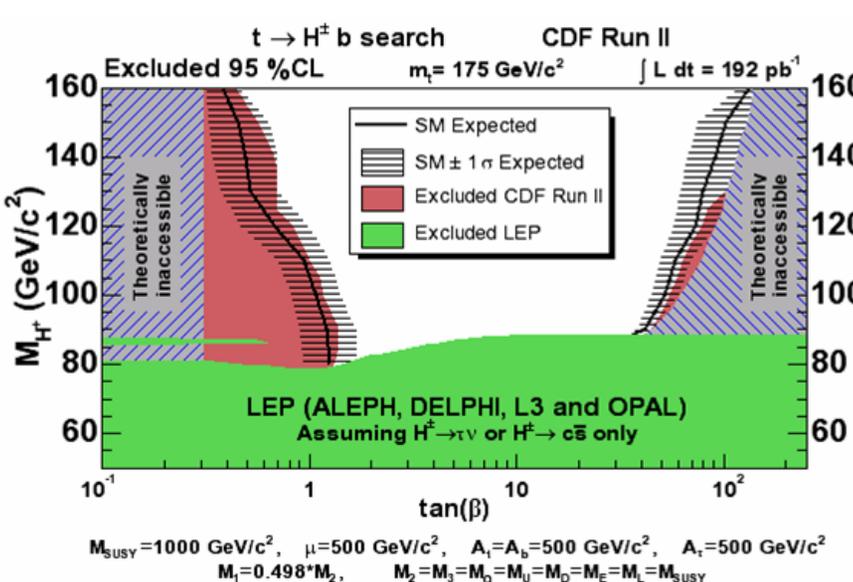


Results: MSSM, Benchmarks 1 to 4

$\mu = -500 \text{ GeV}$

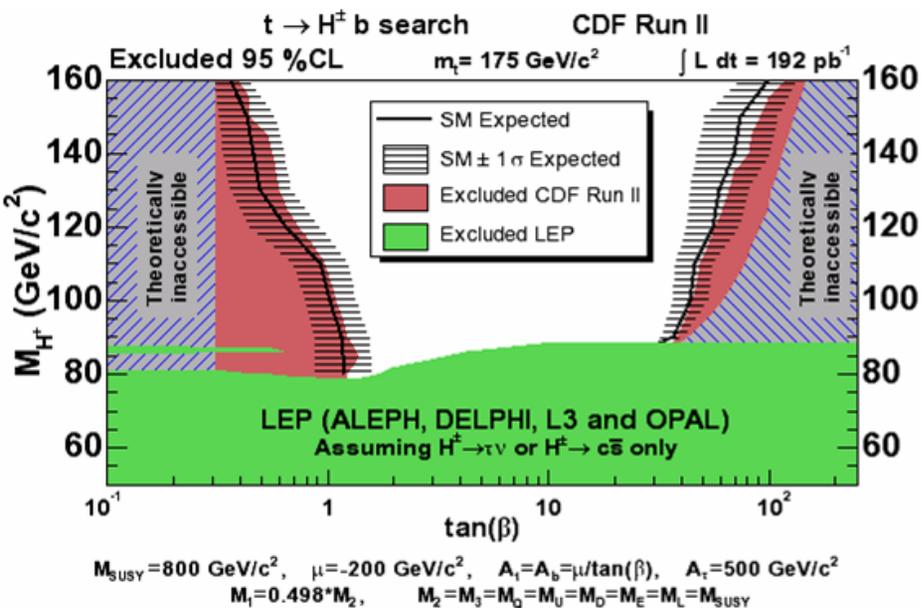


$\mu = 500 \text{ GeV}$

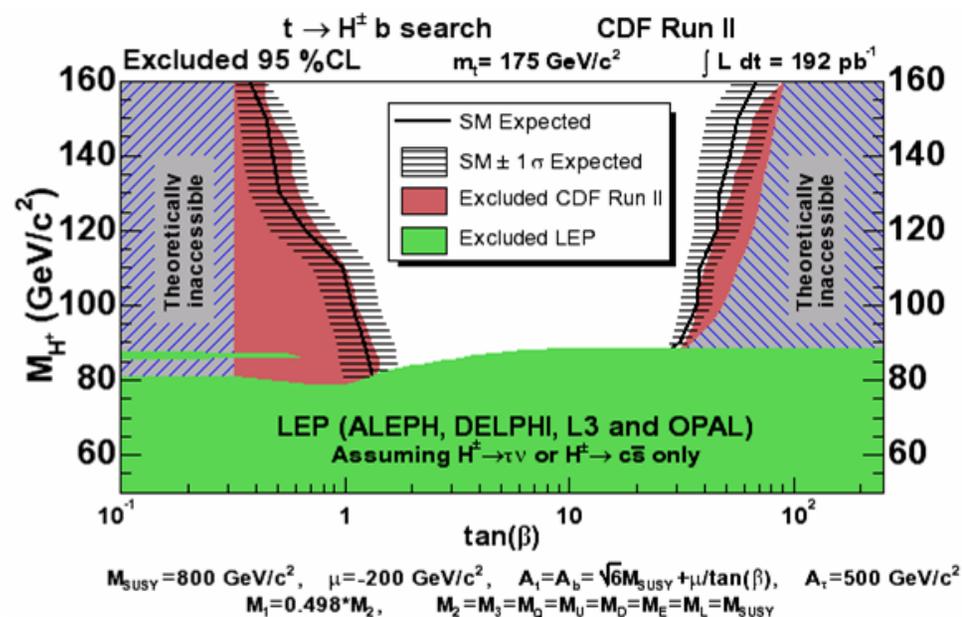


Results: MSSM, Benchmark 5 and 6

Minimal stop Mixing



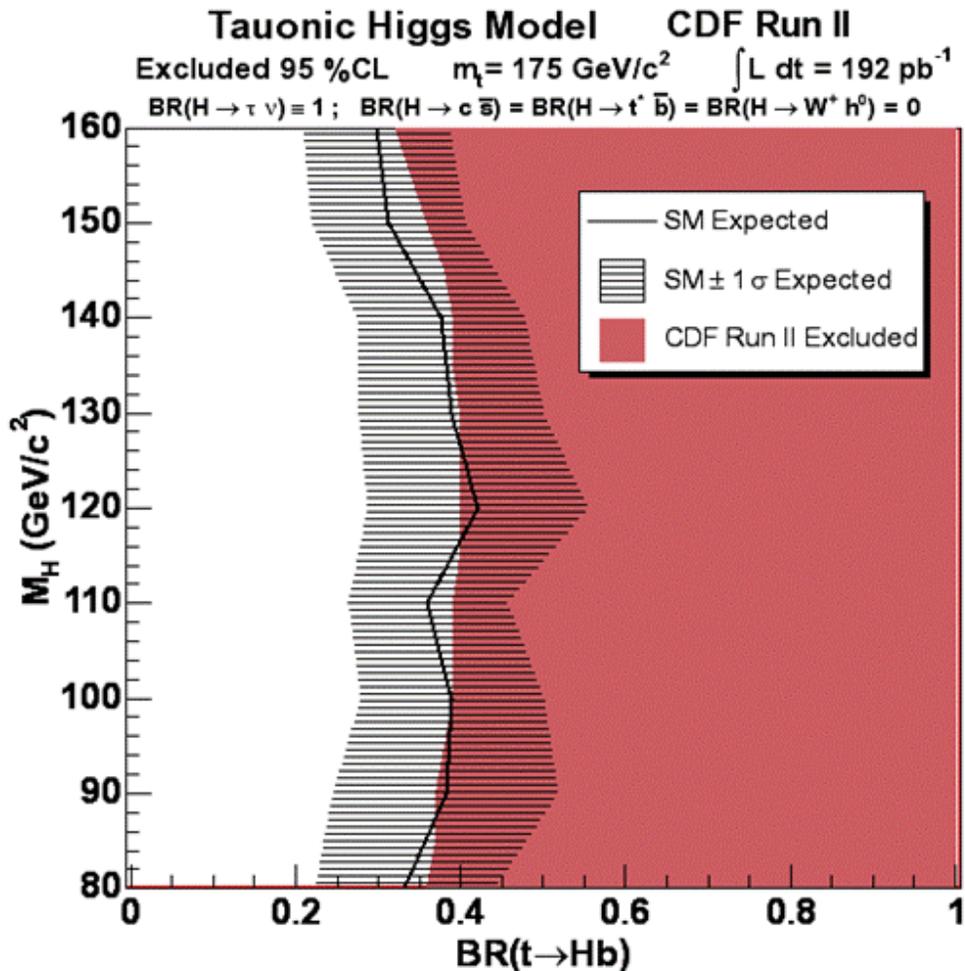
Maximal stop mixing



Clearly, these benchmarks were not designed for the charged Higgs search in $t\bar{t}$ decay products.

Results: Tauonic Higgs Model ($H^+ \rightarrow \tau \nu$)

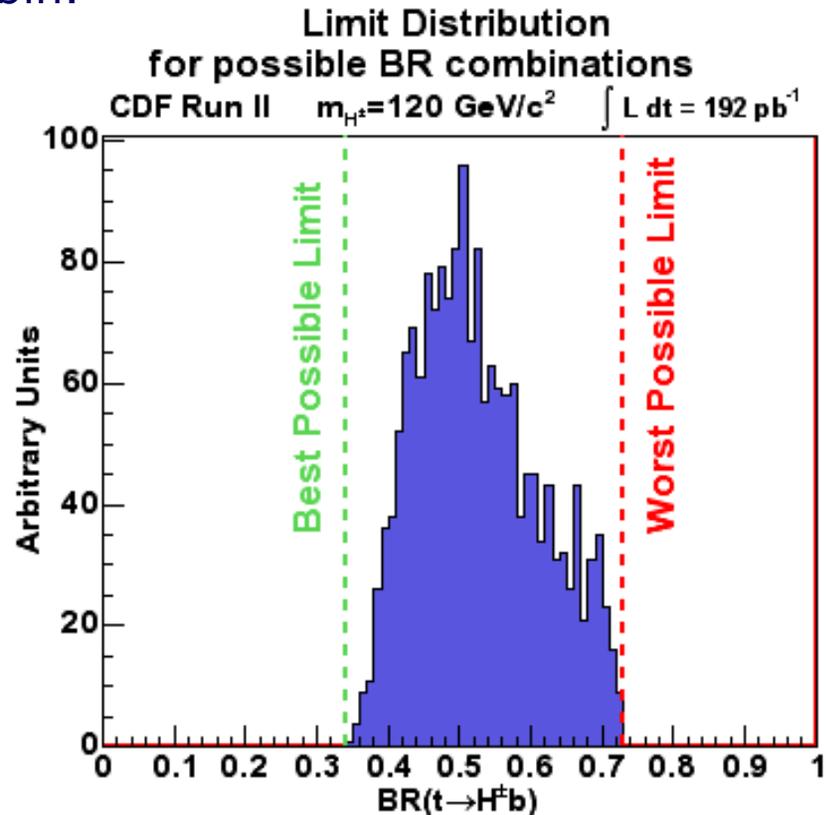
- Set $BR(H^+ \rightarrow cs) \equiv BR(H^+ \rightarrow t^*b) \equiv BR(H^+ \rightarrow Wh^0) \equiv 0$. Limits directly on $BR(t \rightarrow H^+b)$



Results: Higgs BR independent

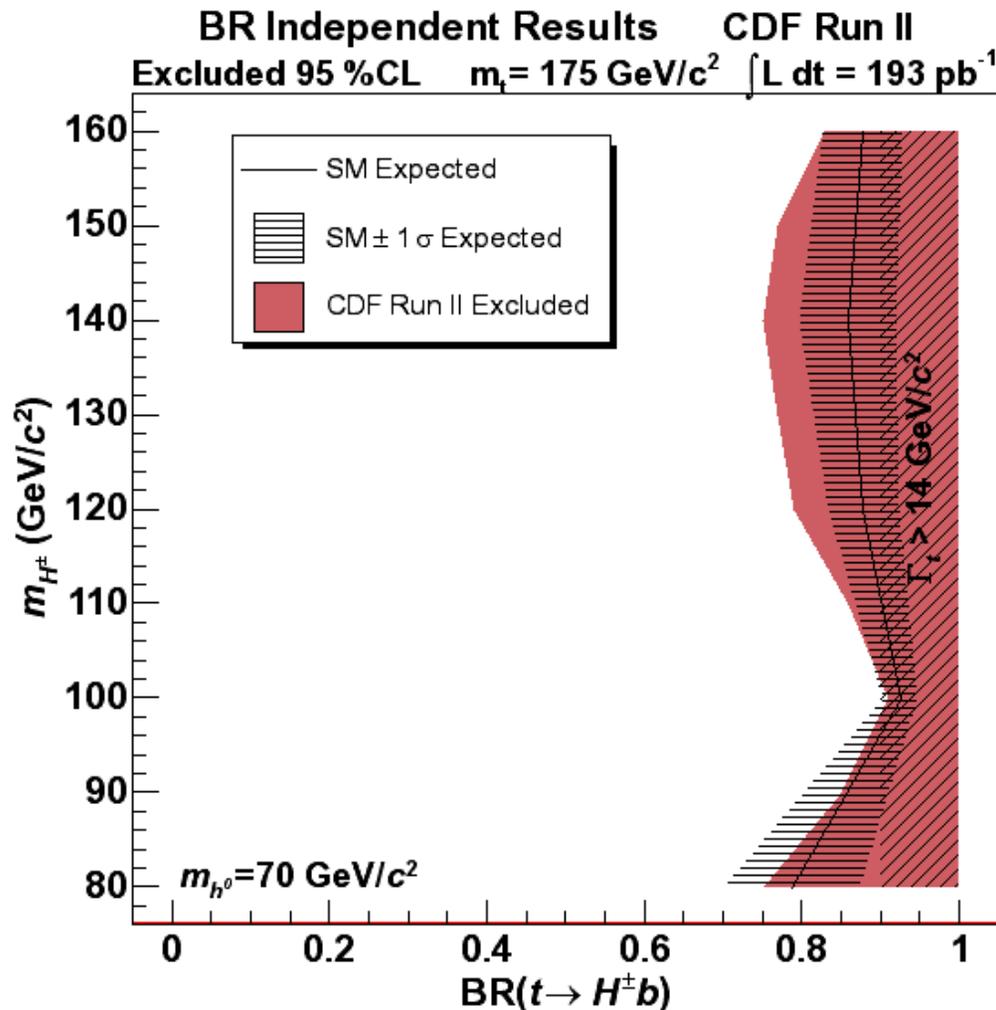
- ➔ **Scan over all BR combinations and take the worst limit:**
 - ➔ Slice the $BR(H^+ \rightarrow cs, t^*b, W^+h^0)$ in bins of 0.05 (21 bins each, 1771 total)
 - ➔ Limit on $BR(t \rightarrow H^+b)$ for each bin.

Limit distribution
for $m_{H^\pm} = 120 \text{ GeV}/c^2$



95%CL limits on $BR(t \rightarrow H^+b)$ ranging from 0.34 to 0.73

Results: Higgs BR independent



BR($t \rightarrow H^\pm b$) < 0.9 @95%CL for 80 GeV < m_{H^\pm} < 160 GeV!

✓ Results obtained in the context of three different models.

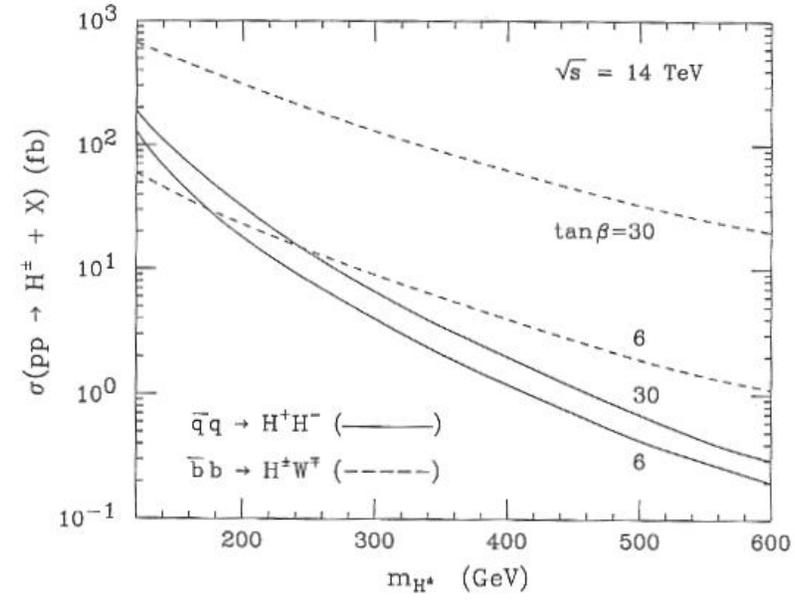
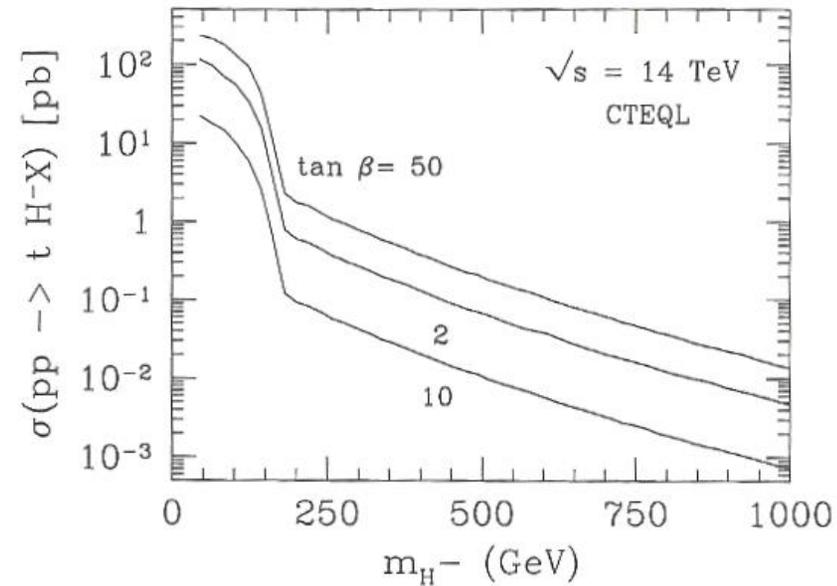
What is expected for the future ?

The Future, LHC

The Future : H^\pm Production at LHC

$$m_{H^+} < m_{Top} - m_b$$

Main source : $t\bar{t}, t \rightarrow H^+ b$
Well behind : $q\bar{q} \rightarrow H^+ H^-$
 $b\bar{b} \rightarrow H^+ W^-$



$$m_{H^+} > m_{Top} - m_b$$

Main source : $gg, q\bar{q} \rightarrow t\bar{b}H^-$
Well behind : $q\bar{q}(gg) \rightarrow H^+ H^-$
 $b\bar{b} \rightarrow H^+ W^-$

*Again, H^\pm stronger signal comes associated with a top quark.
 Strongest signal if $m_H < m_{Top} - m_b$*

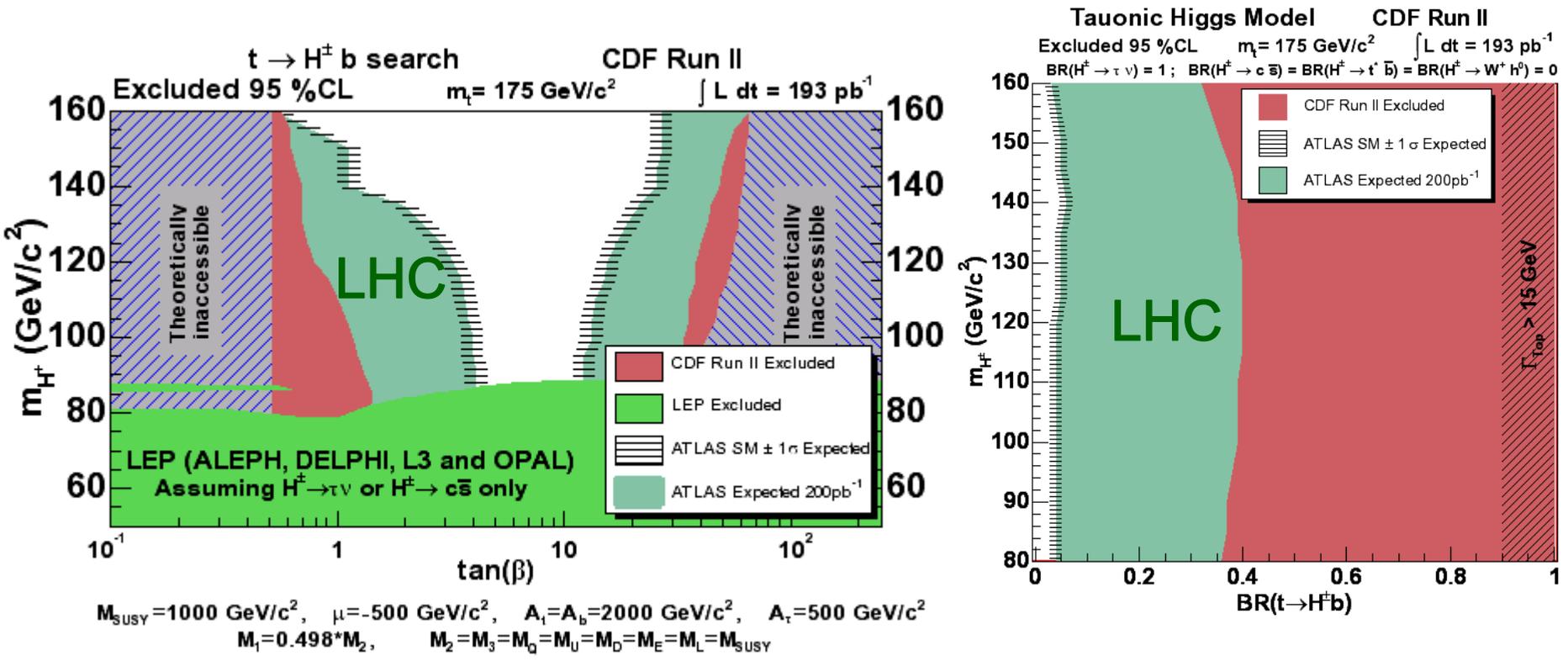
The Future : LHC $t\bar{t}$ Production

- ➔ **From Tevatron : $t\bar{t}$ channels to first order as expected from SM. Either :**
 - ➔ H^\pm does not exist, or it is not between 80 GeV and 160 GeV
 - ➔ the top rarely decays to H^\pm
 - ➔ H^\pm decays are significantly shared between decay modes.
- ➔ **Expect the same at LHC !**
- ➔ **Projections for Atlas(hep-ph/0403021) in 200 pb⁻¹**
 - ➔ Use only 200 pb⁻¹ to compare to the TeV results
 - ➔ At LHC this is only 1 week of nominal luminosity data taking !
 - ➔ Prod. Rate $\sigma_{t\bar{t}} \sim 633$ pb. **wow!**

Channel ($L \equiv e, \mu$)	S/B	# $t\bar{t}$ expected	#Background
Dilepton	10	1600	160
L+Jets (1 ⁺ Tag)	28	5280	185
L+Jets (2 ⁺ Tag)	78 wow-wow!	1740	22
L+ τ_{had}	10	290	29

Use these numbers to calculate raw estimates for the limits on Charged Higgs.

The Future : LHC Limits on Charged Higgs

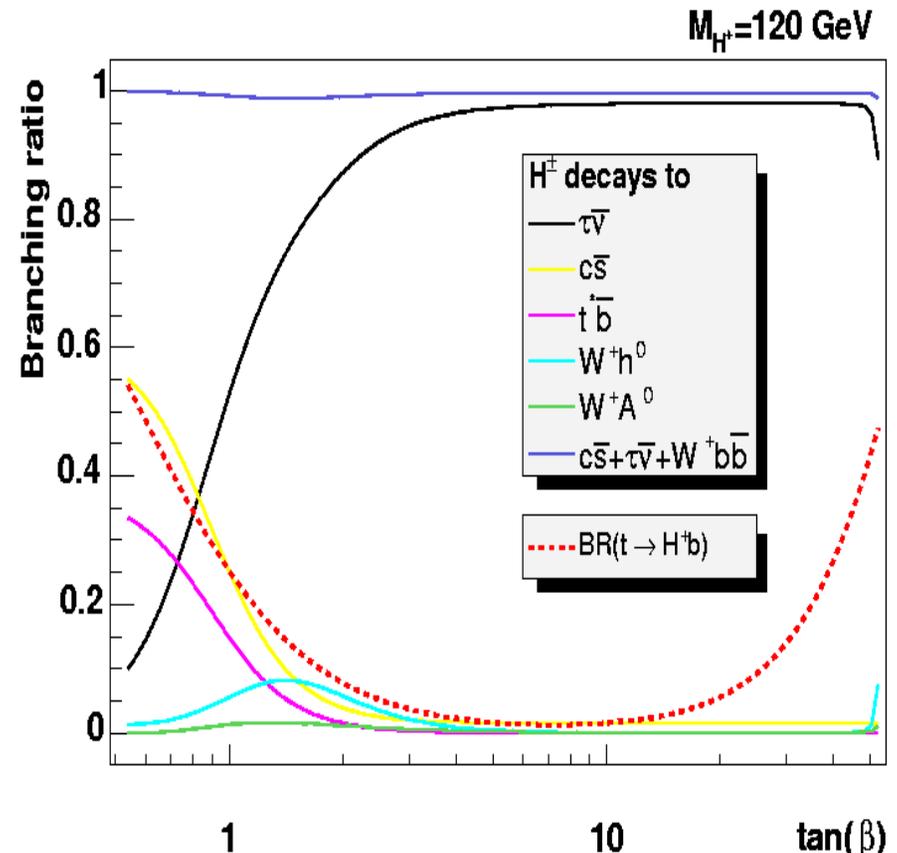
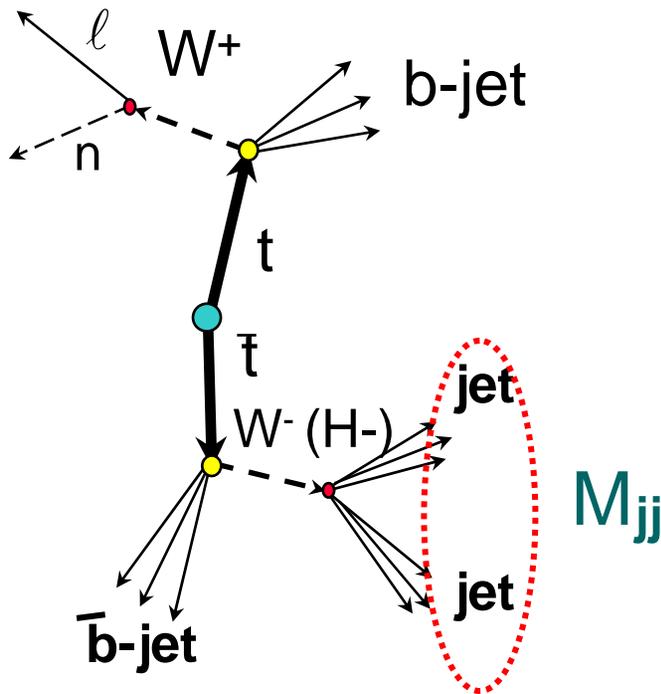


Large exclusion region promptly obtained
Note the small uncertainties

Other ongoing analyses

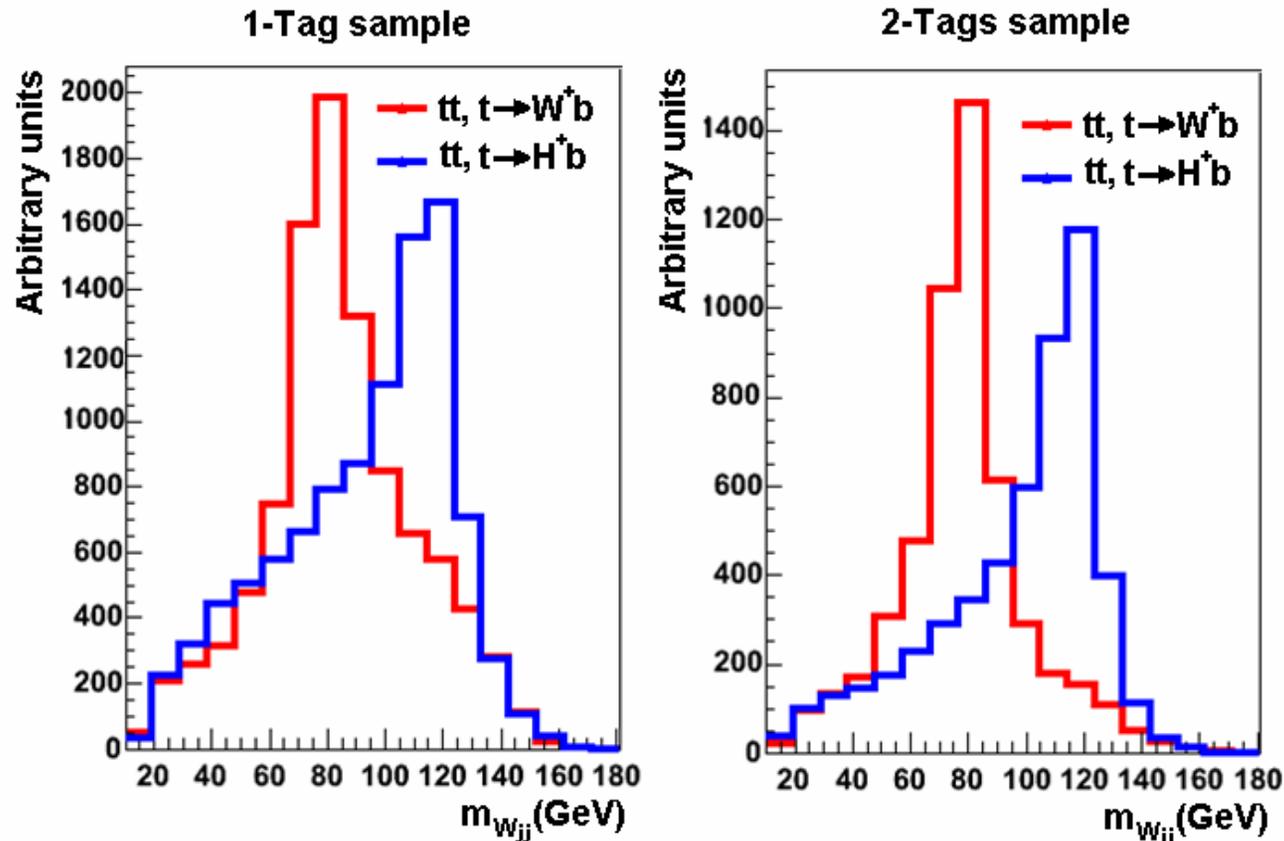
Search for $H^+ \rightarrow c\bar{s}$ in top decays

- ➔ **Strategy** : look for a second bump in W mass distribution
- ➔ Search for $H^+ \rightarrow c\bar{s}$, or generic charged particle decay into two jets
- ➔ In MSSM : $H^+ \rightarrow c\bar{s}$ happens at low $\tan(\beta)$, low m_{H^\pm}



Kinematic Fitter

- Kinematic Fitter: $m_{W_{lv}}=80.4$, $m_t=m_{\bar{t}}=175$ GeV, **$m_{W_{jj}}$ is left unconstrained**
- Charged Higgs Signal Pythia MC: 1M events, $m_t=175$ GeV, $m_{H^+}=120$ GeV/ c^2 , $\tan(\beta)=1$, $BR(H^\pm \rightarrow cs)=1$



Summary

Summary

→ $H^+ \rightarrow \tau \nu$ analysis

- Improved Hadronic tau sensitivity by using jet-tau likelihood fitter.
- Results stricter than previous results using only that channel.

→ $H^+ \rightarrow c s$ analysis

- Brand new. Still to understand shapes for signal, $t\bar{t}$, and W +jets
- Shape of the charged Higgs signals look promising!
- Look for publication early next year.

→ Multi-channel analysis

- Radiative corrections included to the best of our knowledge.
- Benchmark parameters developed specifically for the charged Higgs search.
- **Best limits to date!** See PRL **96**, 042003 (2006)
- Promising results expected at the LHC!

CDF is thoroughly searching for the charged Higgs boson

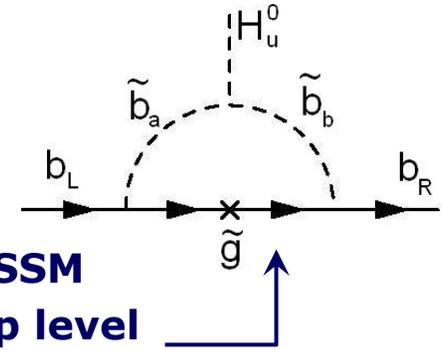
Thank you very much.

Correction to bottom Yukawa coupling

Effective Lagrangian :

$$h_b H_d^0 b \bar{b} + \underbrace{\Delta h_b H_u^0 b \bar{b}}$$

- This term is forbidden at tree-level MSSM
- Δh_b dynamically generated at one-loop level



- Relation between the bottom mass and its Yukawa coupling changes:

$$m_b = v_1 h_b \rightarrow m_b = v_1 h_b + v_2 \Delta h_b = v_1 h_b (1 + \Delta h_b \tan(\beta)) = v_1 h_b (1 + \Delta m_b)$$

- Mass of b well known, this induces a change in effective Yukawa coupling that affects the whole multiplet
Thus, the Yukawa coupling of charged Higgs to top and bottom is modified:
Whenever a m_b is brought from Yukawa coupling replace as follows

$$m_b \rightarrow \frac{m_b}{1 + \Delta m_b}$$

Δm_b = Loops from : gluino-sbottom, charged Higgsino-stop, charged Wino-stop, neutral Wino-sbottom

See Nucl. Phys. B577,88 (2000) for more details.

Number of Expected Events

- Dilepton, lepton+jets $\equiv 1$ and ≥ 2 tags, lepton+tau XS analyses (XSA)

$$\mu_{XSA}^{\text{exp}} = N_{XSA}^{\text{back}} + \sigma \mathcal{E}_{tt, XSA}$$

Includes Luminosity

Number of Expected Events, N^{back}

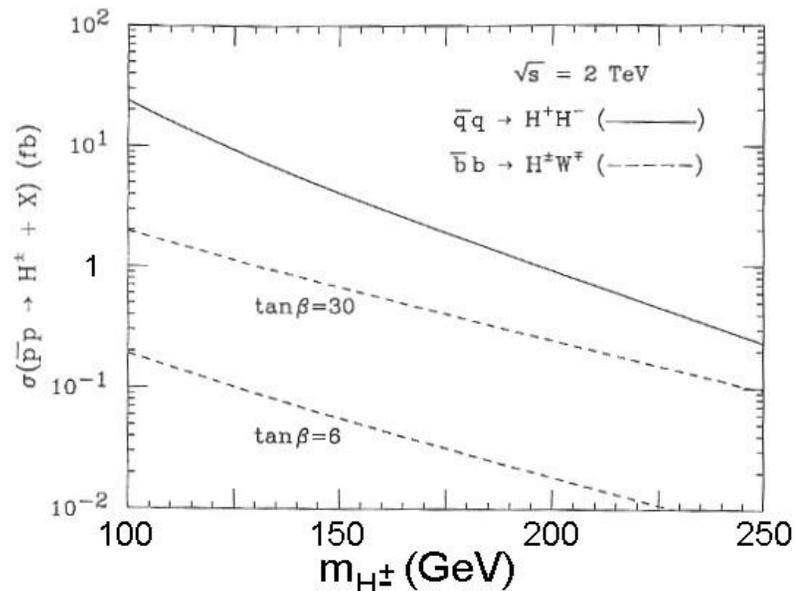
- Dilepton, lepton+jets $\equiv 1$ and ≥ 2 tags, lepton+tau XS analyses (XSA)

$$\mu_{XSA}^{exp} = \underbrace{N_{XSA}^{back}} + \sigma \mathcal{E}_{tt, XSA}$$

Taken from the cross section measurement

- Assume non-SM backgrounds to be negligible

- $pp \rightarrow Wh^0(\text{MSSM}) < \sigma(W h^0_{SM}) < 0.2 \text{ pb}$
- $pp \rightarrow Zh^0(\text{MSSM}) < \sigma(Z h^0_{SM}) < 0.1 \text{ pb}$
- $pp \rightarrow H^+H^-$
- $pp \rightarrow W^+H^-$
- H^+ production via decay of heavy SUSY particles. *Ignored here.*



Number of Expected Events, σ

- ➔ Dilepton, lepton+jets $\equiv 1$ and ≥ 2 tags, lepton+tau XS analyses (XSA)

$$\mu_{XSA}^{\text{exp}} = N_{XSA}^{\text{back}} + \underbrace{\sigma}_{\text{theoretical}} \mathcal{E}_{tt, XSA}$$

Use the theoretical production cross section $\sigma^{\text{theo}} = (6.7 \pm 0.7) \text{pb}$, [hep-ph0303085](#)

- ➔ Assume that introduction of the Higgs sector do not change the production rate.

Number of Expected Events, B_i 's

- Dilepton, lepton+jets $\equiv 1$ and ≥ 2 tags, lepton+tau XSA analyses (XSA)

$$\mu_{XSA}^{\text{exp}} = N_{XSA}^{\text{back}} + \sigma \underbrace{\mathcal{E}_{tt, XSA}}$$

$$\mathcal{E}_{tt, XSA} = \sum_{i,j=1}^5 \underbrace{B_i B_j}_{\substack{\text{B}_i \text{ (B}_j\text{)} : \text{Branching fractions of top} \\ \text{(anti-top) decay mode}}} \mathcal{E}_{i,j \text{ XSA}} \left(\Gamma_{\text{top}}, \Gamma_{\text{Higgs}}, m_{H^\pm}, m_{h^0} \right)$$

Total efficiency calculated from top and anti-top branching ratio decay modes.

- Recall that the B_i 's are calculated assuming the Narrow Width Approximation is valid. The analysis is limited to regions in which the widths of top and Higgs are each below 15 GeV.

Number of Expected Events, $\mathcal{E}_{i,j}^{XSA}$

- ➔ Dilepton, lepton+jets $\equiv 1$ and ≥ 2 tags, lepton+tau XS analyses (XSA)

$$\mu_{XSA}^{\text{exp}} = N_{XSA}^{\text{back}} + \sigma \underbrace{\mathcal{E}_{tt,XSA}}_{\mathcal{E}_{tt,XSA} = \sum_{i,j=1}^5 B_i B_j \underbrace{\mathcal{E}_{i,j}^{XSA}}_{\text{Mode-specific efficiency}} (\Gamma_{\text{top}}, \Gamma_{\text{Higgs}}, m_{H^\pm}, m_{h^0})}$$

Mode-specific efficiency determined given $(\Gamma_{\text{top}}, \Gamma_{\text{Higgs}}, m_{H^\pm}, m_{h^0})$

- ➔ It is the efficiency of the tt event with modes i, j given the mass of the charged Higgs and the mass of the neutral Higgs h^0 .
- ➔ It takes into account corrections due to large width of the top and Higgs.

Present limits

- ⇒ **LEP:** Direct search; $m_{H^\pm} > 78.6 \text{ GeV}$ @ 95 % CL, irrespective of $\tan(\beta)$. Combined ALEPH, DELPHI, L3 and OPAL collaborations.

- ⇒ **CLEO:** Indirect limit; measurement of $b \rightarrow s\gamma$ decay rate results in $m_{H^\pm} > (244 + 63/\tan(\beta))^{1.3} \text{ GeV}$ assuming 2HDM only. Can be circumvented in SUSY.

- ⇒ **Tevatron: Run I, results in the $(m_H, \tan(\beta))$ plane :**
 - ⇒ CDF : Direct search in $t \rightarrow H^+ b \rightarrow \tau \nu b$.
 - ⇒ CDF & D0 : indirect searches using the "Lepton+Jets" (+ "Dilepton" (+ "Dilepton" for CDF) analyses using leading order calculations in similar to this studies.
 - ⇒ D0 : analysis using NN.

Top and Higgs Widths

Widths grow with phase space and couplings :

Couplings :

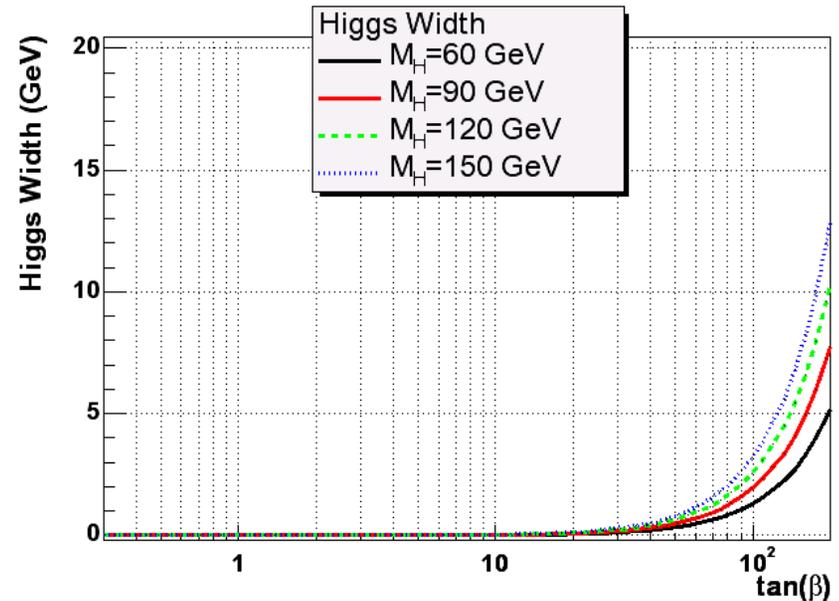
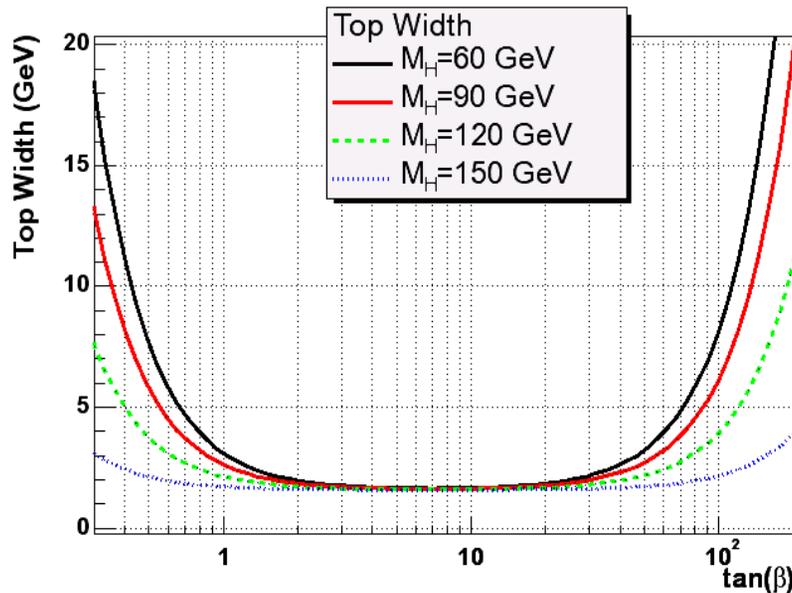
t->Hb coupling get large at high and low tanbeta.

H->τν grows with $\tan(\beta)^2$.

Phase space :

High m_H ; low Γ_{top} high Γ_H

Low m_H ; high Γ_{top} low Γ_H



Higgs width correction, high tanbeta region, accounted for here.

Top width correction accounted too.

Note that it decreases with m_H iggs.

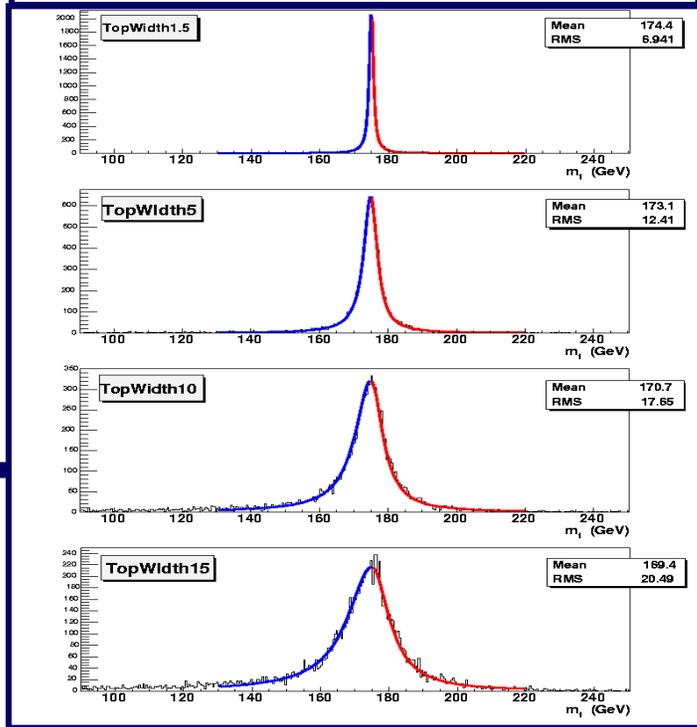
Fairly symmetrical below 15 GeV. **Don't expect and don't see much change in eff.**

Width Corrections, quantitatively

Width correction taken into account when doing :

$$\mathcal{E}_k^{MC}(\Gamma_{top}, \Gamma_{Higgs}, m_{H^\pm}, m_{h^0}) = \int_0^\infty W^t(m'_{top}, \Gamma_{top}) \left(\int_0^\infty \mathcal{E}_k^{rawMC}(m'_{top}, m'_{H^\pm}, m_{h^0}) W^H(m_{H^\pm}, m'_{H^\pm}, \Gamma_{Higgs}) dm'_{H^\pm} \right) dm'_{top}$$

W^t(175, wT) for diff wT



tt → WbWb			
\mathcal{E}_k^{rawMC} wrt $m_{top}=175$	m_{top} (GeV)		
	165	175	185
Dilep	0.94	≡1	1.12
LJets1+	0.96	≡1	1.1
LTauH	0.93	≡1	1.09

After corrections

tt → WbWb				
$\mathcal{E}_k^{MC}(wT, \dots)$ wrt $\Gamma_{top} = 1.4$	Γ_{top} (GeV)			
	1.4	5	10	15
LJets1+	≡1	1.002	1.003	1.005

Width corrections are very small, although we still take them into account.

Removal of Overlap Between XS'

- ➔ Separate the lepton+jets into exactly 1 tag and 2 or more tags
- ➔ Signal Overlap between lepton+jets and dilepton ?
 - ➔ No. Lepton+jets has a dilepton veto
- ➔ Signal Overlap between XSA and lepton+tau, given by
 - ➔ $F_{XSA} \equiv \# \text{ events passing both XSA and lepton+tau} / \# \text{ events passing XSA}$

in (%)	$m_{H^\pm} = 120 \text{ GeV}$		
$t\bar{t} \rightarrow b\bar{b} +$	F_{Dilepton}	$F_{\text{l+jets 1 Tag}}$	$F_{\text{l+jets } \geq 2 \text{ Tags}}$
WW (SM)	0.1±0.1	1.1±0.2	1.5±0.4
$HW \rightarrow \tau\nu W$	0.2±0.2	8.7±0.4	10.1±0.9
$HH \rightarrow \tau\nu\tau\nu$	1.1±0.7	12±1	15±2

~1% for SM

up to 15%
if Higgs is
present

➔ Strategy :

- ➔ Implement a lepton+tau veto cut in the lepton+jets and dilepton XS's.
- ➔ Recalculate signal (and background) efficiencies

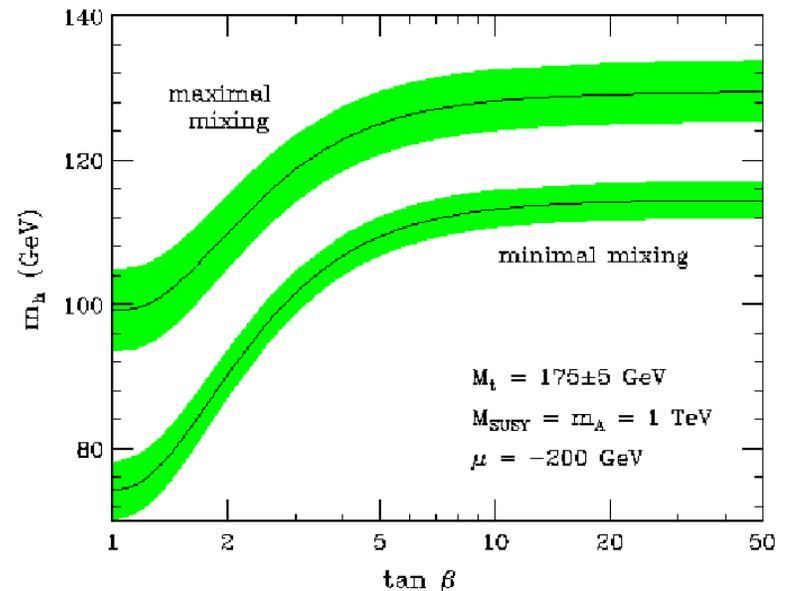
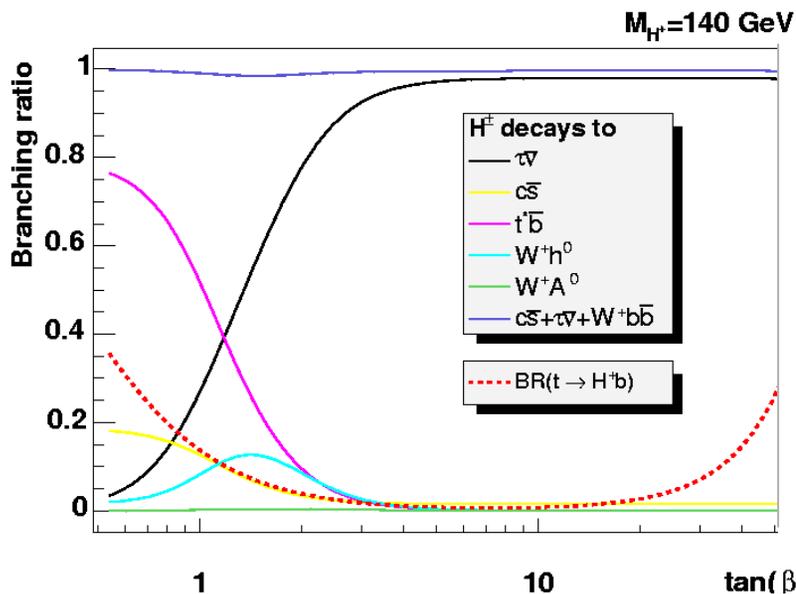
Results: MSSM

- ➔ **Several parameterizations of MSSM**
 - ➔ General MSSM :
 - ➔ intergenerational mixing.
 - ➔ Complex phases.
 - ➔ 105 input parameters in addition to the SM ones.
 - ➔ Phenomenological MSSM :
 - ➔ Soft SUSY breaking parameters are real. No new sources of CP violation.
 - ➔ Matrices for sfermions and trilinear couplings are diagonal. No FCNC at tree level.
 - ➔ Masses and trilinear couplings of 1st and 2nd generation are equal.
 - ➔ 22 input parameters. Down to 14 if only third generation needed.
 - ➔ GUT-constrained MSSM (mSUGRA)
 - ➔ Unification of gaugino masses.
 - ➔ Universal scalar masses.
 - ➔ Universal trilinear coupling.
 - ➔ 4 and a half parameters.
- ➔ **We use the Phenomenological MSSM (pMSSM), with 14 parameters.**
- ➔ **Use different sets of pMSSM parameters (or benchmark scenarios)**

Results: MSSM, Choice of Benchmarks

➔ LEP benchmarks revisited

- ➔ Maximal and minimal stop mixing scenarios.
- ➔ Maximize and minimize the mass of the h^0 as a function of $\tan(\beta)$.
- ➔ All parameters except **At** fixed. **At** is chosen so as to maximize or minimize m_{h^0}

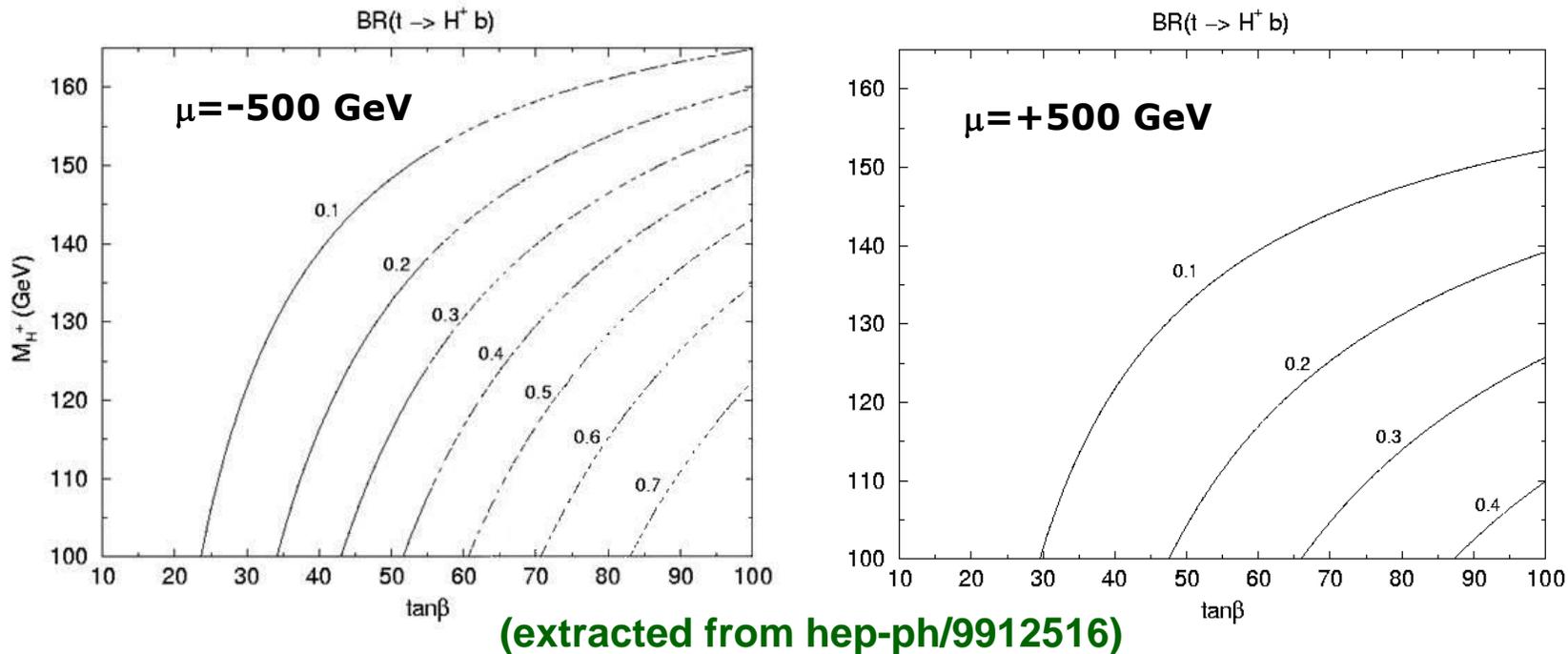


➔ The decay $H^+ \rightarrow W^+ h^0$ larger around $\tan(\beta) \approx 1$

- ➔ $\Gamma(H \rightarrow \tau\nu, c\bar{s}, t^* b)$ is bigger at lower and higher $\tan(\beta)$ values
- ➔ Maximization or minimization of the h^0 mass useful at $\tan(\beta) \approx 1$

Results: MSSM , Choice of Benchmarks

- ➔ **BR($t \rightarrow H^+ b$) strongly depends on the MSSM parameter μ**



- ➔ **Problem** : Previous calculations developed in the large $\tan(b)$ approx.
- ➔ **Recalculated to all ranges of $\tan(b)$. CDF note 7348 (R. Eusebi, M. Carena)**
- ➔ **Summary from the last two slides :**
- ➔ μ has strong effects in the BR($t \rightarrow H^+ b$) predictions at high $\tan(b)$
 - ➔ At has strong effects in BR($H^+ \rightarrow W^+ h^0$) at $\tan(b) \approx 1$.

Results: MSSM , Choice of Benchmarks

	μ (GeV)	A_t (GeV)	Rest of parameters (GeV)
B1	-500	2000	$M_2=M_3=M_Q=M_U=M_D=1$ TeV $M_1=0.4978*M_2, M_L=M_E= 1$ TeV $A_b=A_t, A_{\tau}=500$ GeV
B2	-500	-500	
B3	500	500	
B4	500	2800	
B5(Minimal)	-200	$\mu/\tan(\beta)$	$M_Q=M_U=M_D=1$ TeV, $M_2=M_3=200$ $M_E=M_L=M_Q, A_t=A_b, A_{\tau}=500$
B6(Maximal)	-200	$2450\text{GeV} + \mu/\tan(\beta)$	

- ➔ **B1 and B2 value of $\mu=-500$ GeV, large $BR(t \rightarrow H^+b)$ at large $\tan(b)$**
 - ➔ Difference is A_t , that is chosen so as to maximize (B1) and minimize (B2) the mass of the h^0 in the $\tan(b) \sim 1$ region.
- ➔ **B3 and B4 value of $\mu=+500$ GeV, small $BR(t \rightarrow H^+b)$ at large $\tan(b)$**
 - ➔ Difference is A_t , that is chosen so as to maximize (B4) and minimize (B3) the mass of the h^0 in the $\tan(b) \sim 1$ region.
- ➔ **B5 and B6 are the minimal and maximal stop mixing scenarios used at LEP. They minimize and maximize the mass of the h^0 at every point in $\tan(b)$.**

Results: Higgs BR independent

- ➔ Scan over all BR combinations and take the worst limit:
 - ➔ Slice the Higgs BR($H^+ \rightarrow c\bar{s}$, t^*b , W^+h^0) in bins of 0.05. (21 bins each, 1771 total)
 - ➔ In each bin scan $a \equiv \text{BR}(t \rightarrow H^+b)$ from 0 to 0.9 evaluating the posterior.
 - ➔ In each point in scan : $\text{BR}(h^0 \rightarrow b\bar{b}) \equiv 0.9$, $\Gamma_{\text{Higgs}} = 1 \text{ GeV}$, $\Gamma_{\text{top}} = 1.4/(1 - \text{BR}(t \rightarrow Hb))$

$$P(\alpha, | n_{ll}, n_{lj1}, n_{l\tau}, n_{lj2+}) = \frac{L(n_{ll}, n_{lj1}, n_{l\tau}, n_{lj2+} | \alpha) \pi(\alpha)}{\int L(n_{ll}, n_{lj1}, n_{l\tau}, n_{lj2+} | \alpha') \pi(\alpha') \pi d\alpha'}$$

flat between
0 and 1

- ➔ Integrate the posterior and obtain the 95% CL
- ➔ Repeat for all bins and take worst limit.
- ➔ Repeat for different charged Higgs masses

Results in the ($m_H, \text{BR}(t \rightarrow Hb)$) plane