



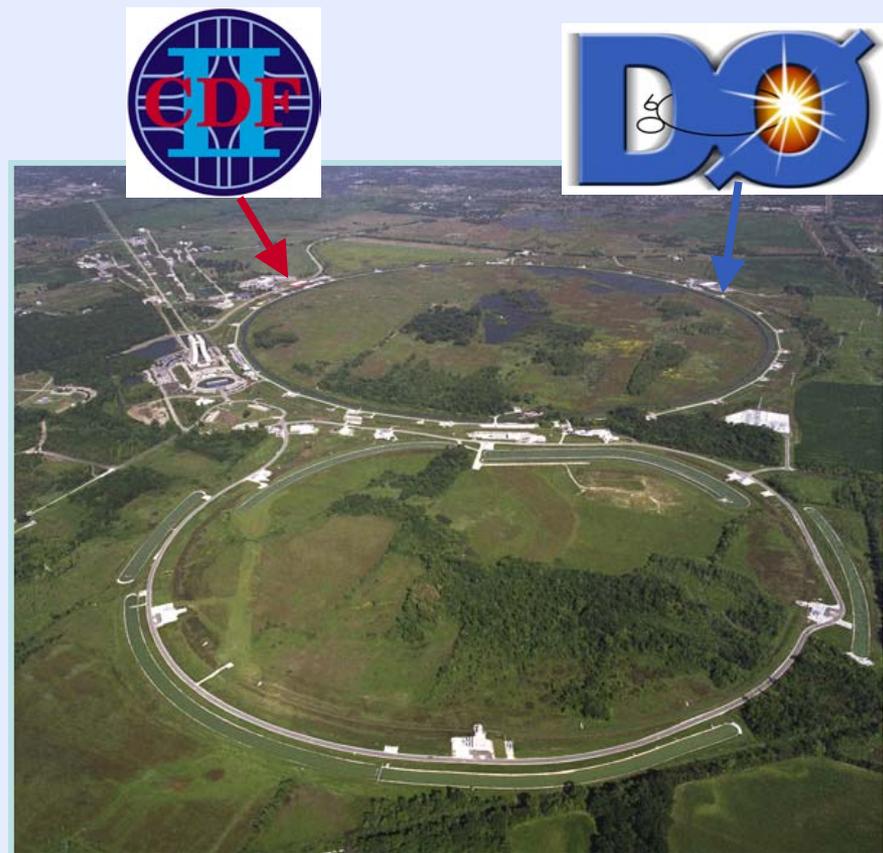
CP violation in B hadrons at the Tevatron

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for the CDF and DØ collaborations

*Recontres de Moriond QCD, March 2010
La Thuile, Italy*



- CDF and DØ detectors at Fermilab
- CP violation in B_s decays
 - Neutral B_s system
 - Mixing induced CP violation
 - Phenomenology
 - Fit results
 - Combined results for CP violating parameter β_s
 - Preview of coming updates:
 - B flavour tagging improvements
 - inclusion of $B_s \rightarrow J/\psi KK$ (S-wave)
 - PID extension
- Summary



The Fermilab Tevatron

CP violation in neutral B_s system

Flavour eigenstates:

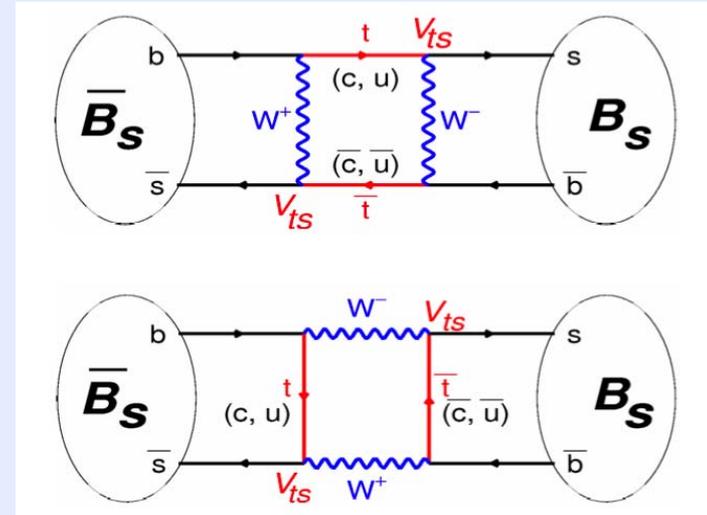
$$\begin{aligned} |B_s^0\rangle &= (\bar{b}s) \\ |\bar{B}_s^0\rangle &= (b\bar{s}) \end{aligned}$$

Mixing of flavour eigenstates is governed by:

$$i \frac{d}{dt} \begin{pmatrix} B_s^0(t) \\ \bar{B}_s^0(t) \end{pmatrix} = H \begin{pmatrix} B_s^0(t) \\ \bar{B}_s^0(t) \end{pmatrix} \equiv \left[\underbrace{\begin{pmatrix} M_0 & M_{12} \\ M_{12}^* & M_0 \end{pmatrix}}_{\text{mass matrix}} - \frac{i}{2} \underbrace{\begin{pmatrix} \Gamma_0 & \Gamma_{12} \\ \Gamma_{12}^* & \Gamma_0 \end{pmatrix}}_{\text{decay matrix}} \right] \begin{pmatrix} B_s^0(t) \\ \bar{B}_s^0(t) \end{pmatrix}$$

Flavour eigenstates are not mass eigenstates:

$$\begin{aligned} |B_s^H\rangle &= p |B_s^0\rangle - q |\bar{B}_s^0\rangle \\ |B_s^L\rangle &= p |B_s^0\rangle + q |\bar{B}_s^0\rangle \end{aligned}$$



Different masses -> mixing frequency:

$$\Delta m_s = m_H - m_L \approx 2|M_{12}|$$

-> phase:

$$\varphi_s^{\text{SM}} = \arg(-M_{12}/\Gamma_{12}) \sim 0.04$$

Different decay widths:

$$\Delta\Gamma = \Gamma_L - \Gamma_H \approx 2|\Gamma_{12}| \cos(2\varphi_s^{\text{SM}})$$

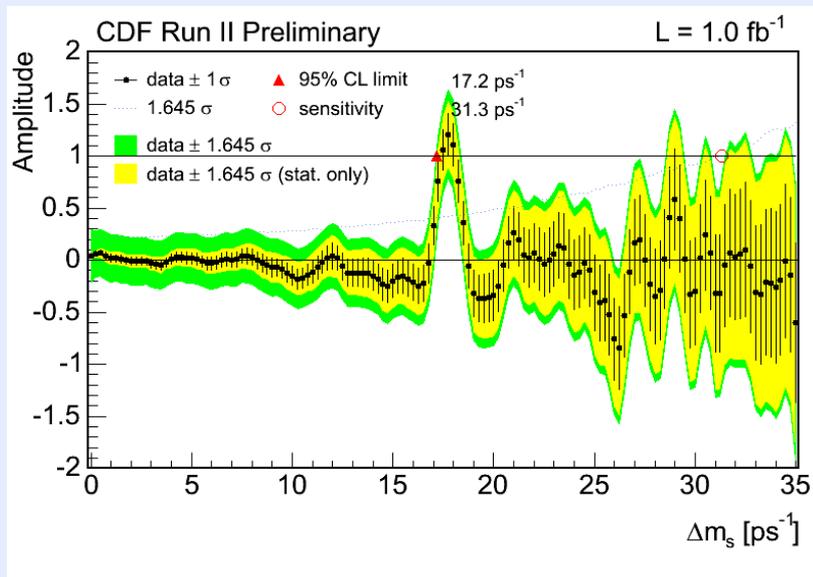


CP violation in $B_s \rightarrow J/\psi\phi$

Δm_s is already well determined by Tevatron measurements

CDF

DØ

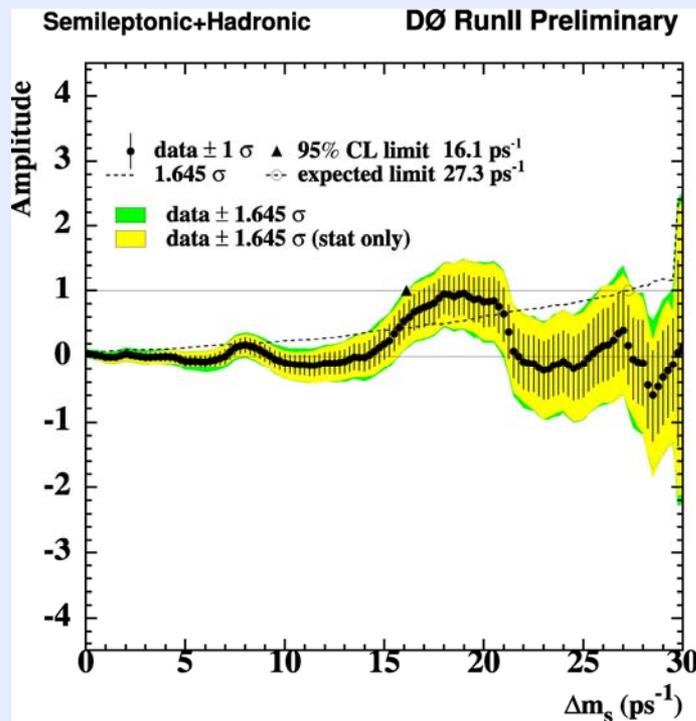


5 σ observation

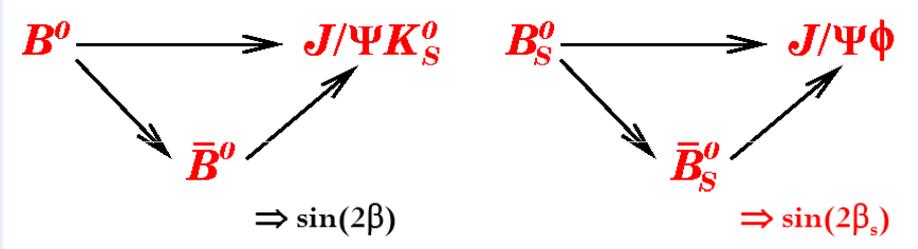
$$\Delta m_s = 17.77 \pm 0.10(\text{stat}) \pm 0.07(\text{stat})ps^{-1}$$

3 σ evidence - in agreement with CDF

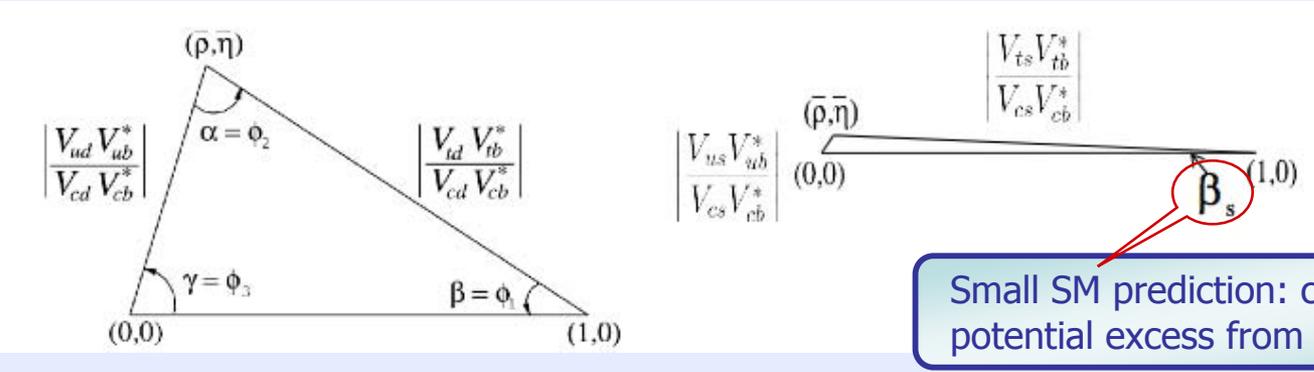
$$\Delta m_s = 18.52 \pm 0.91ps^{-1}$$



However, CP violation phase is not precisely known...
both are needed to constrain New Physics



CP violation in $B_s \rightarrow J/\psi\phi$ occurs through interference of decays with and without mixing.



$$\beta_s = \beta' \equiv \arg \left(-\frac{V_{ts} V_{tb}^*}{V_{cs} V_{cb}^*} \right)$$

Small SM prediction: clear to see potential excess from NP

A New Physics effect would contribute to both the phases φ_s and β_s by introducing a new physics phase:

$$\varphi_s = \varphi_s^{\text{SM}} + \varphi_s^{\text{NP}} \text{ and } 2\beta_s = 2\beta_s^{\text{SM}} - \varphi_s^{\text{NP}}$$

So, if NP phase dominates we measure $2\beta_s \approx -\varphi_s \approx \varphi_s^{\text{NP}}$

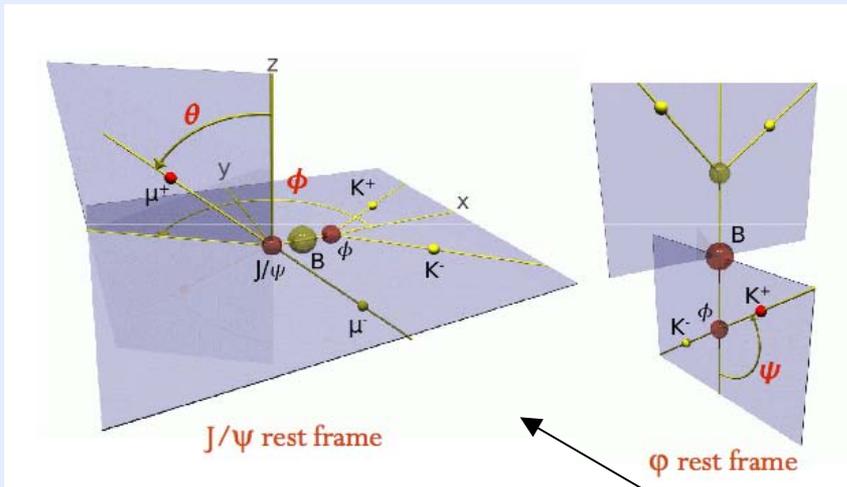
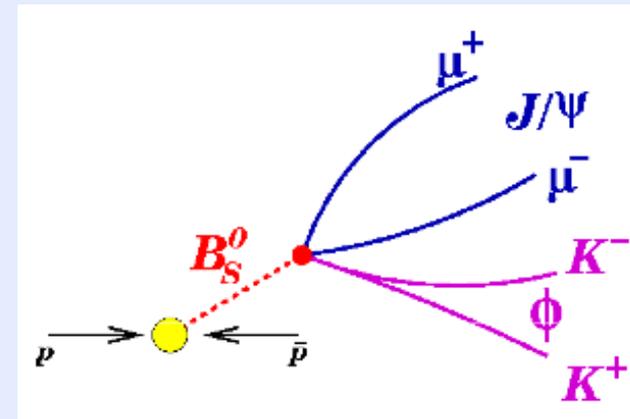
Can measure this in the CP asymmetry:

$$\mathcal{A}_{CP}(t) \equiv \frac{\bar{B}_s^0(t) - B_s^0(t)}{\bar{B}_s^0(t) + B_s^0(t)} = \sin(2\beta_s) \cdot \sin \Delta m_s t$$

$B_s \rightarrow J/\psi \phi$: experimental strategy

Reconstruct $B_s \rightarrow J/\psi (\rightarrow \mu^+ \mu^-) \phi (\rightarrow K^+ K^-)$

Final state is a mixture of CP even ($\sim 75\%$) and odd ($\sim 25\%$) states.



Three angular momentum states of $J/\psi \phi$:

$L=0$	S-wave	CP even
$L=1$	P-wave	CP odd
$L=2$	D-wave	CP even

Can separate final CP states using angular analysis

Transversity basis describes these contributions as: A_0 , $A_{//}$ (CP even), A_{perp} (CP odd) according to their polarisation.

Can be separated using the angular distributions of the final state particles

Use a multivariate fit combining angular analysis and time dependence

- Simplest case, fit without flavour tagging, has four fold ambiguity:
 - β_s and $\Delta\Gamma$ symmetric about zero
 - strong phases symmetric about π

$$\phi_{\parallel} \equiv |A_{\parallel}^* A_0|$$

$$\phi_{\perp} \equiv |A_{\perp}^* A_0|$$

$$\beta_s \rightarrow \frac{\pi}{2} - \beta_s$$

$$\Delta\Gamma \rightarrow -\Delta\Gamma$$

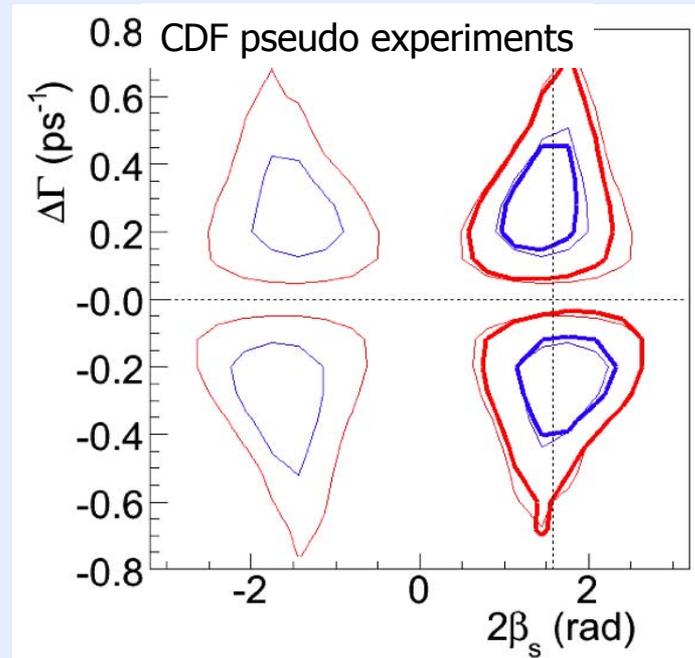
$$\phi_{\parallel} \rightarrow 2\pi - \phi_{\parallel}$$

$$\phi_{\perp} \rightarrow \pi - \phi_{\perp}$$

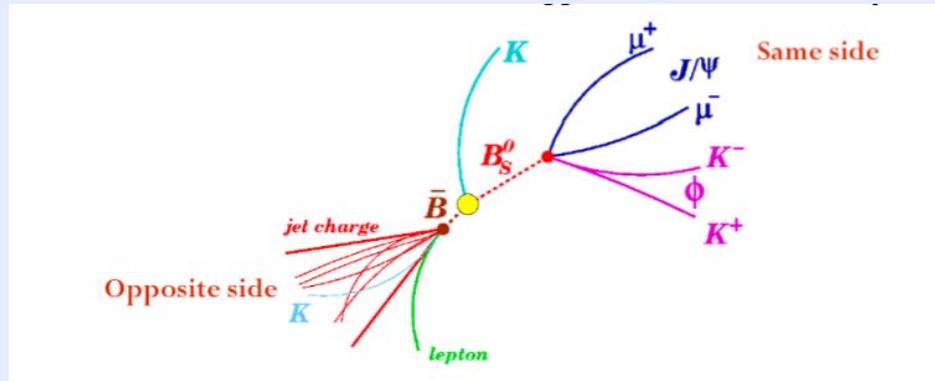
- Addition of flavour tagging allows us to follow time dependence of B_s and \bar{B}_s separately
- Removes insensitivity to sign of $\Delta\Gamma$ and β_s
-> Removes half of the ambiguity

Plot shows CDF pseudo experiments

- 95% CL tagged fit
- 68% CL tagged fit
- 95% CL untagged fit
- 68% CL untagged fit



- Opposite side tag (OST):
 - b quarks are pair produced (flavour conservation)
 - Can deduce properties of the candidate B meson by looking at the decay of the B hadron formed by the pair produced partner of its b quark
- Same side kaon tag (SSKT):
 - Sign of kaon from primary vertex of candidate B_s can tag B_s flavour



- CDF 2.8fb^{-1} result does not exploit SSKT beyond 1.4fb^{-1}
- Updated SSKT now validated and will be included in upcoming result
 - Significant improvement in tagging power
 - Reduces statistical errors of $\beta_s - \Delta\Gamma$ contours

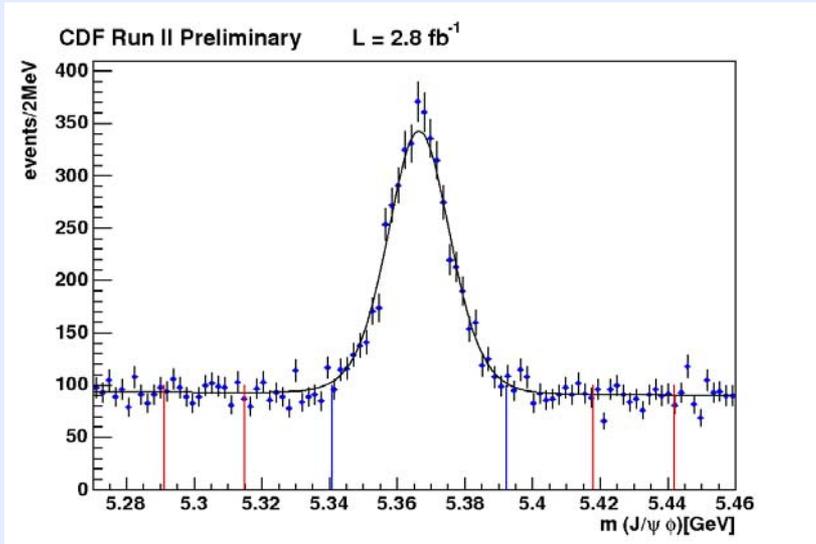
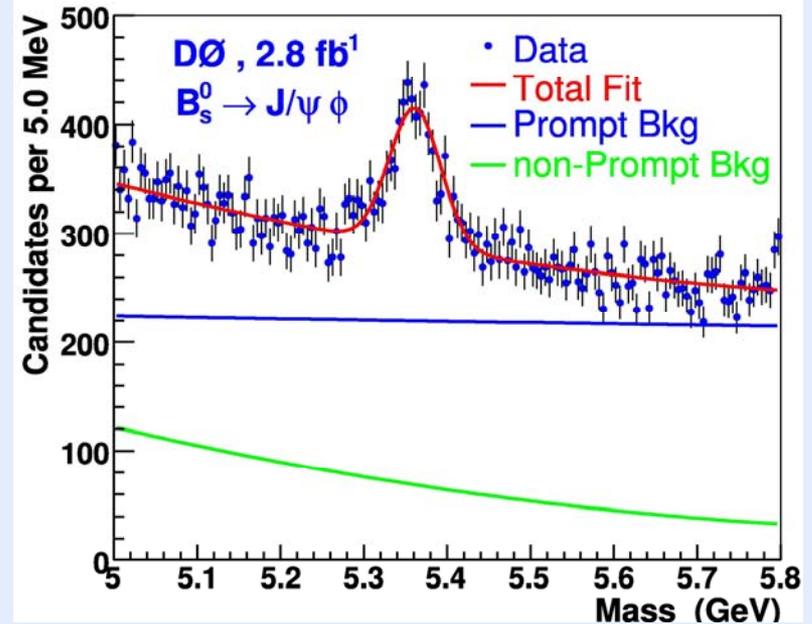
CP violation in $B_s \rightarrow J/\psi\phi$:

- ▣ *CDF and DØ results $L=2.8\text{fb}^{-1}$*
- ▣ *Tevatron combination*



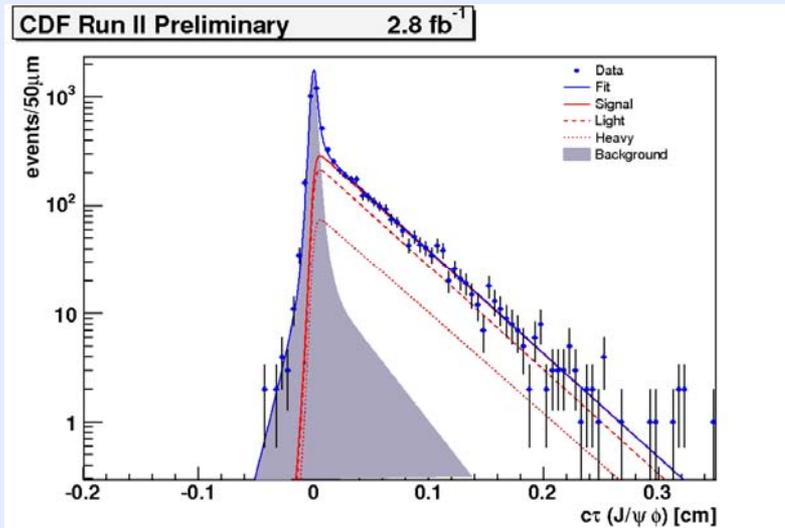
DØ

- Dimuon trigger
- Cut based selection
- Signal events: 1967 ± 65
- Luminosity: 2.8fb^{-1}



CDF

- Dimuon trigger
- Neural network selection
- Signal events: 3153 ± 55
- Luminosity: 2.8fb^{-1}

DØ L=2.8fb⁻¹CDF L=2.8fb⁻¹<http://www-d0.fnal.gov/Run2Physics/WWW/results/prelim/B/B58/>

$$\tau_s = 1.53 \pm 0.04 \text{ (stat.)} \pm 0.01 \text{ (syst.)} ps$$

$$\Delta\Gamma = 0.02 \pm 0.05 \text{ (stat.)} \pm 0.01 \text{ (syst.)} ps^{-1}$$

$$|A_{\parallel}|^2 = 0.241 \pm 0.019 \text{ (stat.)} \pm 0.007 \text{ (syst.)}$$

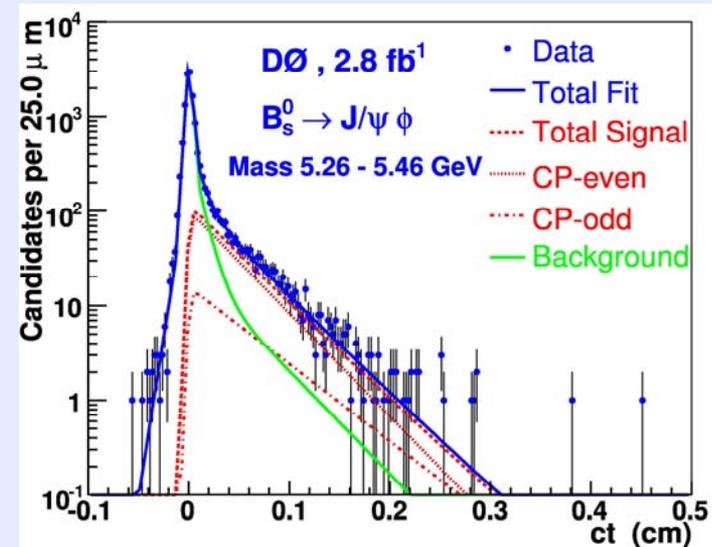
$$|A_0|^2 = 0.508 \pm 0.024 \text{ (stat.)} \pm 0.008 \text{ (syst.)}$$

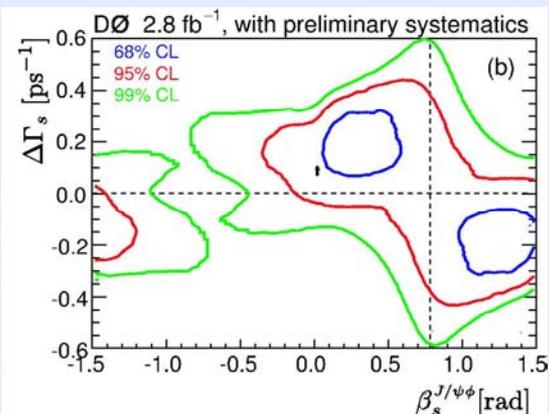
$$\tau_s = 1.52 \pm 0.05 \text{ (stat.)} \pm 0.01 \text{ (syst.)} ps$$

$$\Delta\Gamma = 0.19 \pm 0.07 \text{ (stat.)} {}^{+0.02}_{-0.01} \text{ (syst.)} ps^{-1}$$

$$|A_{\parallel}|^2 = 0.244 \pm 0.032 \text{ (stat.)} \pm 0.014 \text{ (syst.)}$$

$$|A_0|^2 = 0.555 \pm 0.027 \text{ (stat.)} \pm 0.006 \text{ (syst.)}$$

From untagged fit with $\beta_s = 0.0$ http://www-cdf.fnal.gov/physics/new/bottom/080724.blessed-tagged_BsJPsiPhi_update_prelim



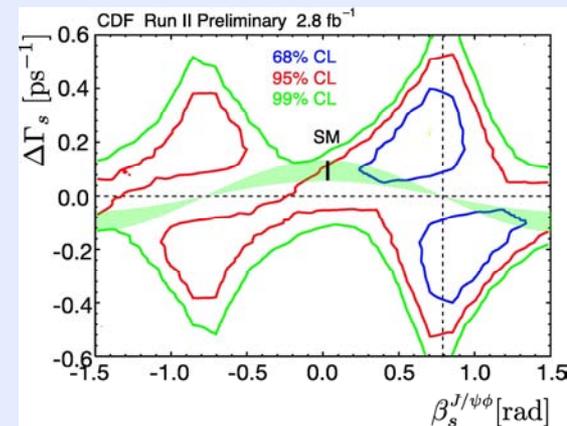
Cannot quote point value for β_s

□ symmetries in likelihood

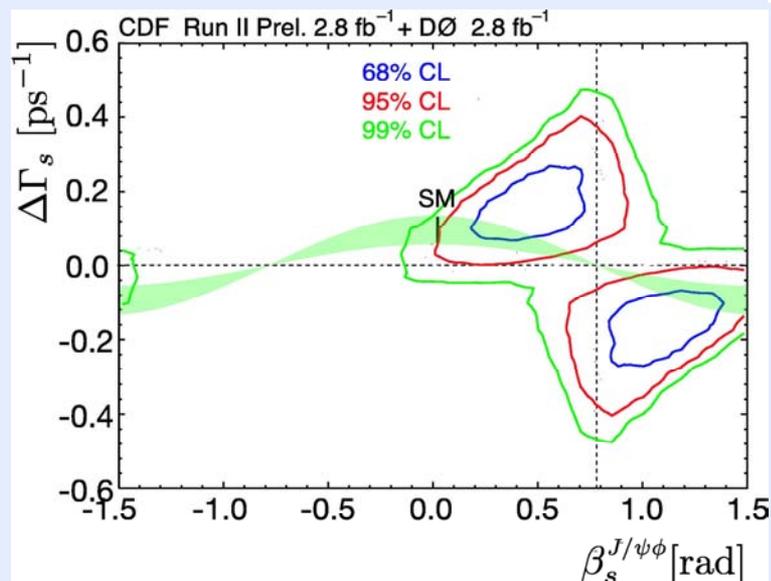
□ non-Gaussian errors

Show $\Delta\Gamma_s$ - β_s likelihood contours

□ errors adjusted by likelihood ratio ordering technique



DØ: 2.8fb^{-1} result
significance 1.7σ of
deviation from SM



CDF: 2.8fb^{-1} result
P-value for SM
point = 7%
→ significance 1.8σ

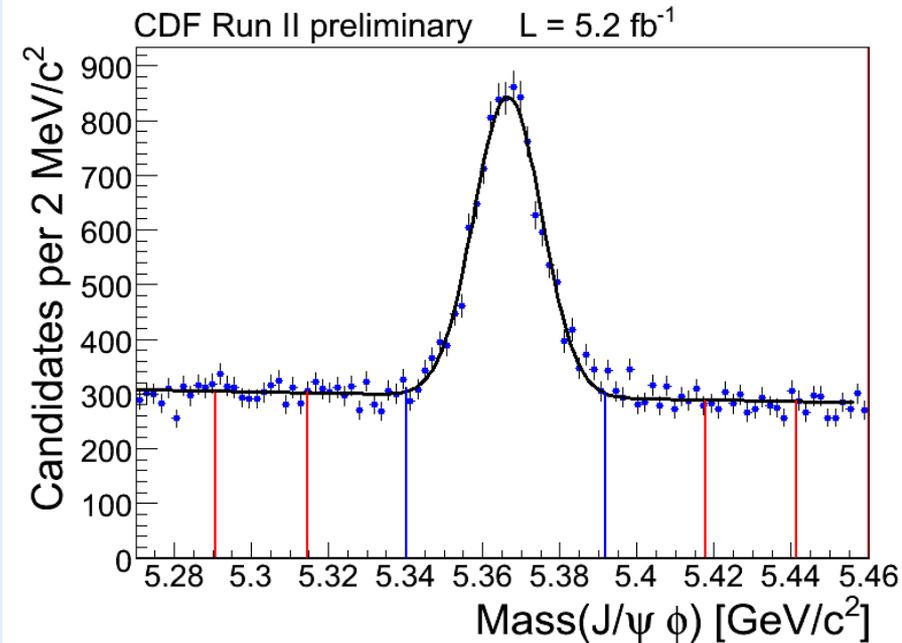
Tevatron combination: probability of observed deviation from SM = 3.4% (2.12σ)

http://tevbwg.fnal.gov/results/Summer2009_betas/

*Coming upgrades to β_s
measurements*

DØ:

- Update to $L=6.1\text{fb}^{-1}$ from previous 2.8fb^{-1} result
- Will use boosted decision tree selection to improve signal to background ratio



CDF:

- Update to $L=5.2\text{fb}^{-1}$
- Inclusion of PID for full dataset
 - previous update lacked PID after 1.4fb^{-1}
 - dE/dx and TOF calibrated and extended
- Improved flavour tagging: addition of same side tag
- Account for $B_s^- \rightarrow J/\psi KK$ S-wave

- SSKT is calibrated on B_s mixing measurement
- Mixing amplitude ≈ 1 :
 - tagger assesses its performance accurately
- Amplitude > 1
 - tagger underestimates its power
- Amplitude < 1
 - tagger overestimates performance
- Uses several decay modes:

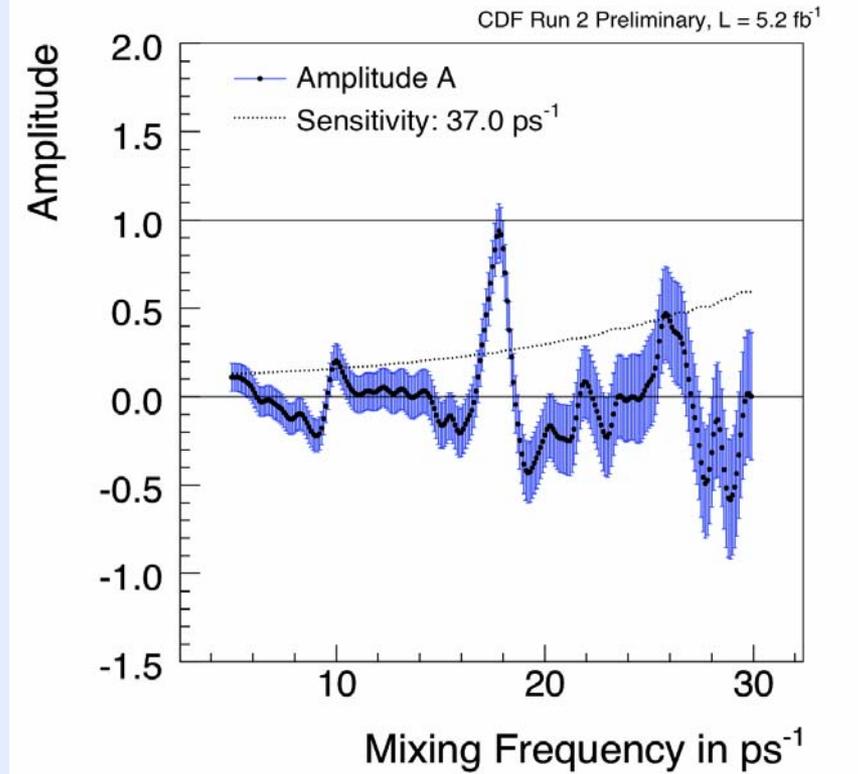
$$B_s^0 \rightarrow D_s^- \pi^+, D_s^- \rightarrow \phi^0 \pi^-, \phi^0 \rightarrow K^+ K^-$$

$$B_s^0 \rightarrow D_s^- \pi^+, D_s^- \rightarrow K^* K^-, K^* \rightarrow K^+ \pi^-$$

$$B_s^0 \rightarrow D_s^- \pi^+, D_s^- \rightarrow (3\pi)^-$$

$$B_s^0 \rightarrow D_s^- (3\pi)^+, D_s^- \rightarrow \phi^0 \pi^-, \phi^0 \rightarrow K^+ K^-$$

Agreement between this and the published CDF measurement is very good

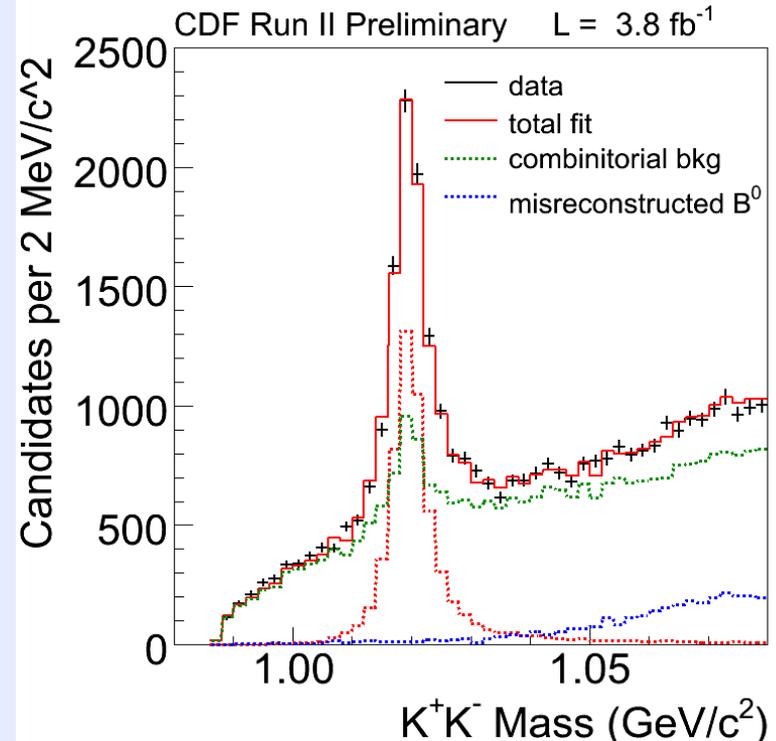
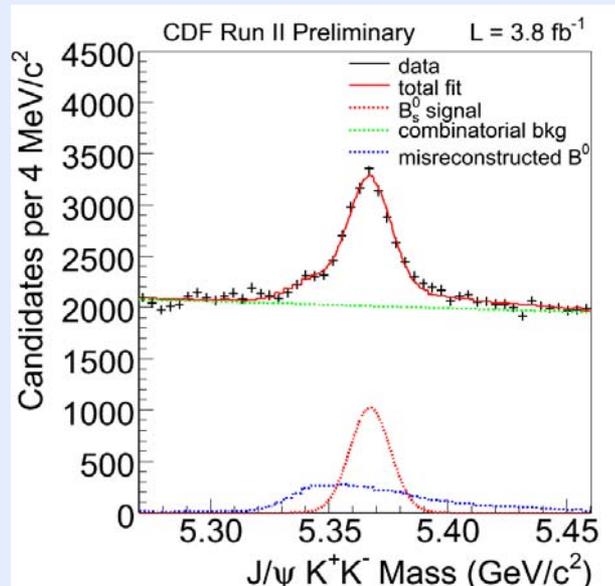


$$\mathcal{A} = 0.94 \pm 0.15 \text{ (stat.)} \pm 0.13 \text{ (syst.)}$$

$$\Delta m_s = 17.79 \pm 0.07 \text{ ps}^{-1}$$

$$\epsilon \mathcal{A}^2 D^2 \approx 3.2 \pm 1.4 \%$$

- Potential contamination of B_s signal by S-wave: $B_s \rightarrow J/\psi KK$ (KK non-resonant) and $B_s \rightarrow J/\psi f^0$
- Predicted up to 10-15% contamination could bias towards SM value of β_s
- Next CDF update will include non-resonant component in the likelihood

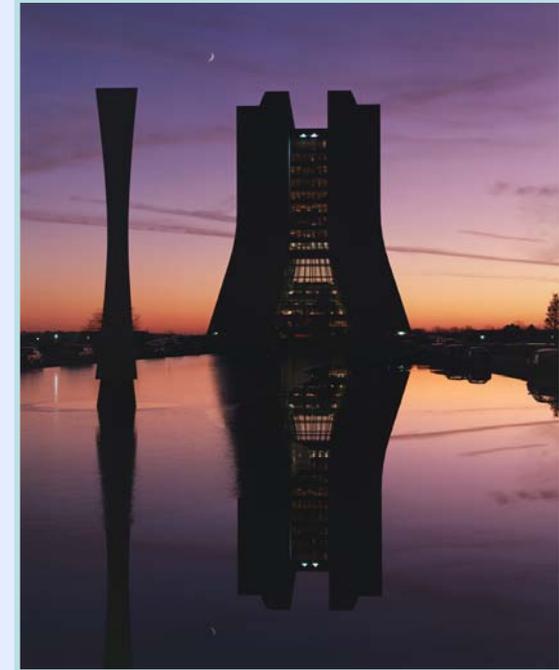


Invariant KK mass (above)

- combinatorial background from B_s sidebands
- B^0 reflections modelled from MC
- Fractions fixed from B_s mass fit (left)

-> little evidence from this distribution for additional component

- B_s system is rich source of information about CP violation, **accessible to Tevatron experiments**
- DØ and CDF produced precision measurements of B_s mixing frequency
- Both experiments have constrained CP violating phase β_s
 - More powerful result achieved from Tevatron combination
- Forthcoming updates from CDF and DØ to utilize full datasets doubling current statistics



More news on this important measurement soon!

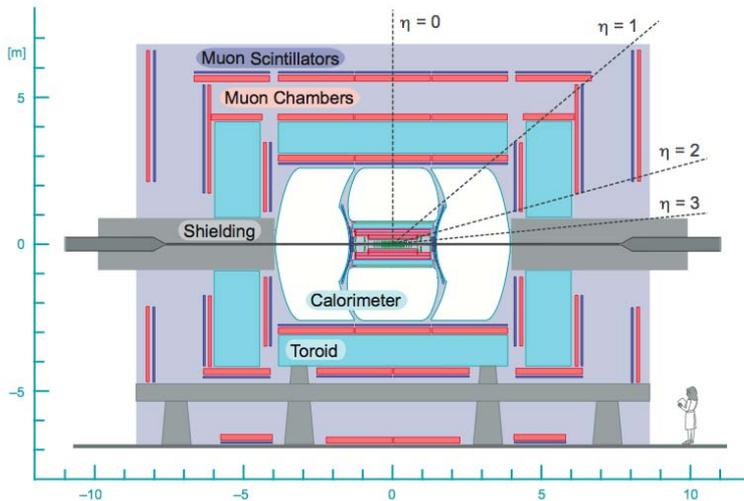
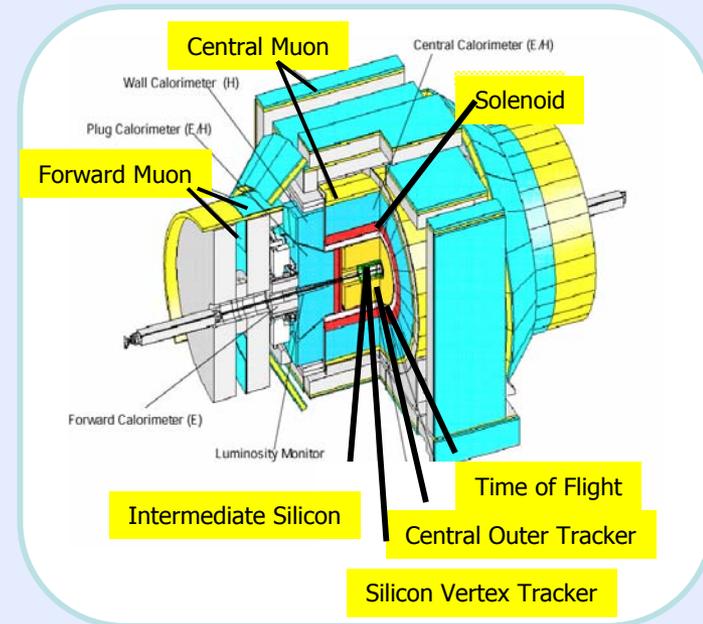
Back up

B physics at CDF:

- Particle ID: dE/dx and TOF
- Excellent vertex resolution $\sim 23\mu\text{m}$ and p_T resolution: $\sigma(p_T)/p_T^2 \sim 0.1 (\text{GeV}/c)^{-1}$
- Trigger level silicon tracking

B physics at DØ:

- Solenoid (2TeV) polarity reversed weekly
- Strengths in semileptonic and J/ψ decays
- Excellent calorimetry and electron ID

**Hadron colliders vs B factories:**

- + Much larger B production cross section, phase space, range of Bs generated
- Higher background, don't know initial state
- > Larger signal for B_s at hadron machines but need sophisticated trigger and selection

Fit function

Use a multivariate fit combining angular analysis and time dependence

- Simplest case: fit without flavour tagging:

$$\mathcal{L}_i = f_s \cdot \underline{P_s(m)} \cdot \underline{T(t, \psi, \theta, \phi)} \cdot \underline{P_s(\sigma_t)} + (1 - f_s) \cdot \underline{P_b(m)} \cdot \underline{P_b(t, \sigma_t)} \cdot \underline{P_b(\sigma_t)} \cdot \underline{P_b(\psi)} \cdot \underline{P_b(\theta)} \cdot \underline{P_b(\phi)}$$

signal

background

mass terms

time dependence and angular terms

lifetime error terms

- Untagged likelihood has four fold ambiguity:

- β_s and $\Delta\Gamma$ symmetric about zero
- strong phases symmetric about π

$$\phi_{\parallel} \equiv |A_{\parallel}^* A_0|$$

$$\phi_{\perp} \equiv |A_{\perp}^* A_0|$$

$$\begin{aligned} \beta' &\rightarrow \frac{\pi}{2} - \beta', \\ \Delta\Gamma &\rightarrow -\Delta\Gamma, \\ \phi_{\parallel} &\rightarrow 2\pi - \phi_{\parallel}, \\ \phi_{\perp} &\rightarrow \pi - \phi_{\perp}. \end{aligned}$$

- Addition of flavour tagging allows us to follow time dependence of B_s and \bar{B}_s separately
- Removes insensitivity to sign of $\Delta\Gamma$ and β_s
- > Removes half of the ambiguity

$$\mathcal{L}_i = f_s \cdot P_s(m) \cdot \underline{P_s(\xi)} \cdot \underline{T(t, \psi, \theta, \phi, \mathcal{D}, \xi)} \cdot P_s(\sigma_t) \cdot \underline{P_s(\mathcal{D})} + (1 - f_s) \cdot P_b(m) \cdot \underline{P_b(\xi)} \cdot \underline{P_b(t, \sigma_t)} \cdot P_b(\psi) \cdot \underline{P_b(\theta)} \cdot \underline{P_b(\phi)} \cdot P_b(\sigma_t) \cdot \underline{P_b(\mathcal{D})}$$

terms altered or added by tag decision or tagging dilution

Use same likelihood ratio ordering technique to account for non-Gaussian behaviour (to ensure we do not undercover the confidence regions) and to include effect of systematics on the errors:

- Generate pseudo experiments at the SM point in the $\Delta\Gamma$ - β_s plane.
- Fit with all parameters floating
- Fit again with $\Delta\Gamma$ and β_s fixed to the SM point
- Form a likelihood ratio:

$$\mathcal{LR} = 2 \log \frac{\mathcal{L}(\beta_s^{J/\psi\phi}, \Delta\Gamma, \vec{\xi})}{\mathcal{L}(\vec{\xi})}$$

- By integrating and normalising the distribution of likelihood ratios, and taking the log of this distribution, get the distribution (1-confidence level)
- This distribution can be used to adjust the likelihood to correspond to that expected from Gaussian errors, for a given confidence level.
- For systematic adjustment of the errors, pseudo experiments are generated for an ensemble of 16 “alternative universes”, for each the nuisance parameters are generated within +/- 5sigma of the measured value. The “universe” with the most extreme variations is used the adjustment at each point.

Polarisation of vector mesons w.r.t direction of motion:

$|A_0|^2$: polarisation longitudinal, parallel

$|A_{//}|^2$: polarisation transverse, parallel

$|A_{\text{perp}}|^2$: polarisation transverse, perpendicular

We let the A 's be normalized such that $|A_0|^2 + |A_{//}|^2 + |A_{\perp}|^2 = 1$.

The predicted angular distributions can be found from the following prescription Let \hat{n} be the unit vector in the direction of the l^+ (J/ψ rest frame),

$$\hat{n} = (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta),$$

and let A be a complex vector defined as

$$\mathbf{A} = (A_0 \cos \psi, -\frac{A_{//} \sin \psi}{\sqrt{2}}, i \frac{A_{\perp} \sin \psi}{\sqrt{2}}).$$

The angular distributions are governed by the probability density

$$P(\theta, \phi, \psi) = \frac{9}{16\pi} |\mathbf{A} \times \hat{n}|^2.$$

This is normalized such that

$$\int \int \int \frac{9}{16\pi} |\mathbf{A} \times \hat{n}|^2 \sin \theta d\theta d\phi \sin \psi d\psi = 1.$$